Nonlinear Motion

Sextupoles cause nonlinear dynamics, which can be chaotic and unstable.

$$\begin{pmatrix} x_{n+1} \\ x'_{n+1} \end{pmatrix} = \underline{M}_0 \left[\begin{pmatrix} x_n \\ x'_n \end{pmatrix} - \frac{k_2 l_s}{2} \begin{pmatrix} 0 \\ x_n^2 \end{pmatrix} \right] \qquad \begin{pmatrix} x_{n+1} \\ x'_{n+1} \end{pmatrix} = \begin{pmatrix} \sqrt{\beta} & 0 \\ -\frac{\alpha}{\sqrt{\beta}} & \frac{1}{\sqrt{\beta}} \end{pmatrix} \begin{pmatrix} \hat{x}_n \\ \hat{x}'_n \end{pmatrix}$$

$$\begin{pmatrix} \hat{x}_{n+1} \\ \hat{x}'_{n+1} \end{pmatrix} = \begin{pmatrix} \cos\mu & \sin\mu \\ -\sin\mu & \cos\mu \end{pmatrix} \left[\begin{pmatrix} \hat{x}_n \\ \hat{x}'_n \end{pmatrix} - \frac{k_2 l_s}{2} \sqrt{\beta} \begin{pmatrix} 0 \\ \beta \hat{x}_n^2 \end{pmatrix} \right]$$

$$\begin{pmatrix} \hat{x}_f \\ \hat{x}'_f \end{pmatrix} = \frac{k_2 l_s}{2} \beta^{\frac{3}{2}} \begin{pmatrix} 1-\cos\mu & \sin\mu \\ -\sin\mu & 1-\cos\mu \end{pmatrix}^{-1} \begin{pmatrix} 0 \\ \hat{x}_f^2 \end{pmatrix} = \frac{k_2 l_s}{2} \beta^{\frac{3}{2}} \frac{1}{2\sin\frac{\mu}{2}} \begin{pmatrix} -\cos\frac{\mu}{2} \\ \sin\frac{\mu}{2} \end{pmatrix} \hat{x}_f^2$$

$$\left.\begin{aligned} \hat{x}_f &= -\frac{4}{k_2 l_s} \beta^{-\frac{3}{2}} \tan\frac{\mu}{2} \\ \hat{x}'_f &= \frac{4}{k_2 l_s} \beta^{-\frac{3}{2}} \tan^2\frac{\mu}{2} \end{aligned}\right\} \hat{x} = \hat{x}_f + \Delta\hat{x} \qquad J_f = \frac{1}{2}(\hat{x}_f^2 + \hat{x}'^2_f) = \frac{1}{2\beta^3}\left(\frac{4}{k_2 l_s} \frac{\tan\frac{\mu}{2}}{\cos\frac{\mu}{2}}\right)^2$$

$$\begin{pmatrix} \Delta\hat{x}_{n+1} \\ \Delta\hat{x}'_{n+1} \end{pmatrix} = \begin{pmatrix} \cos\mu & \sin\mu \\ -\sin\mu & \cos\mu \end{pmatrix} \left[\begin{pmatrix} \Delta\hat{x}_n \\ \Delta\hat{x}'_n \end{pmatrix} - \begin{pmatrix} 0 \\ \frac{k_2 l_s}{2} \beta^{\frac{3}{2}} \Delta\hat{x}_n^2 - 4\tan\frac{\mu}{2} \Delta\hat{x}_n \end{pmatrix} \right]$$

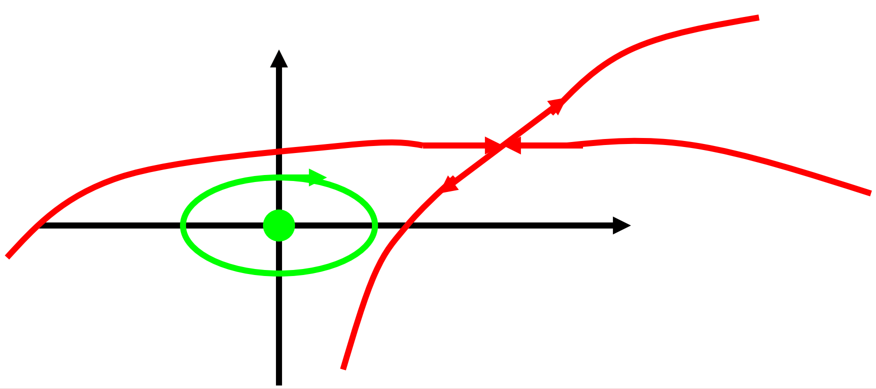
The Dynamic Aperture

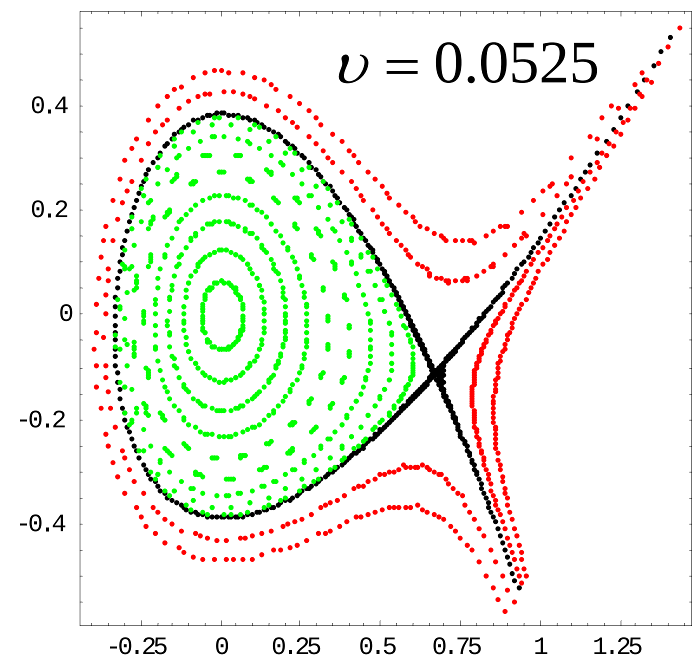
$$\begin{pmatrix}\Delta\hat{x}_{n+1}\\ \Delta\hat{x}'_{n+1}\end{pmatrix}=\begin{pmatrix}\cos\mu & \sin\mu\\ -\sin\mu & \cos\mu\end{pmatrix}\left[\begin{pmatrix}\Delta\hat{x}_{n}\\ \Delta\hat{x}'_{n}\end{pmatrix}-\begin{pmatrix}0\\ \frac{k_2 l_s}{2}\beta^{\frac{3}{2}}\Delta\hat{x}_n^2-4\tan\frac{\mu}{2}\Delta\hat{x}_n\end{pmatrix}\right]$$

$$\begin{pmatrix}\Delta\hat{x}_{n+1}\\ \Delta\hat{x}'_{n+1}\end{pmatrix}=\begin{pmatrix}\cos\mu+4\sin\mu\tan\frac{\mu}{2} & \sin\mu\\ -\sin\mu+4\cos\mu\tan\frac{\mu}{2} & \cos\mu\end{pmatrix}\left[\begin{pmatrix}\Delta\hat{x}_{n}\\ \Delta\hat{x}'_{n}\end{pmatrix}-\begin{pmatrix}0\\ \frac{k_2 l_s}{2}\beta^{\frac{3}{2}}\Delta\hat{x}_n^2\end{pmatrix}\right]$$

$$Tr[\underline{M}]=2\frac{\cos\frac{\mu}{2}(1+2\sin^2\frac{\mu}{2})}{\cos\frac{\mu}{2}}\geq 2$$

The additional fixed point is unstable !





Sextupole Aperture

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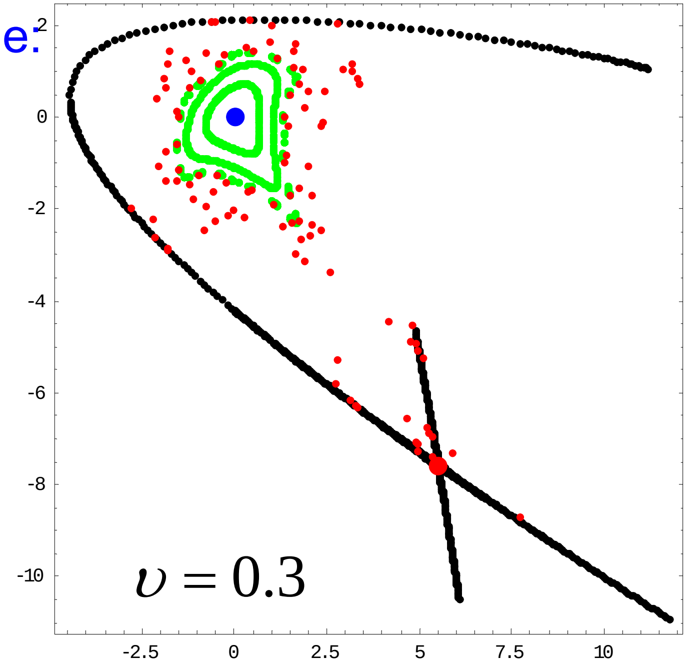
If the chormaticity is corrected by a single sextupole:

$$\xi_x = \xi_{0x} + \frac{1}{4\pi}\beta_x\eta_x k_2 l \approx 0$$

$$J_f = \frac{1}{2\beta^3}\left(\frac{4}{k_2 l_s}\frac{\tan\frac{\mu}{2}}{\cos\frac{\mu}{2}}\right)^2 \approx \frac{1}{2\beta}\left(\frac{\eta}{\xi_0\pi}\frac{\sin\frac{\mu}{2}}{\cos^2\frac{\mu}{2}}\right)^2$$



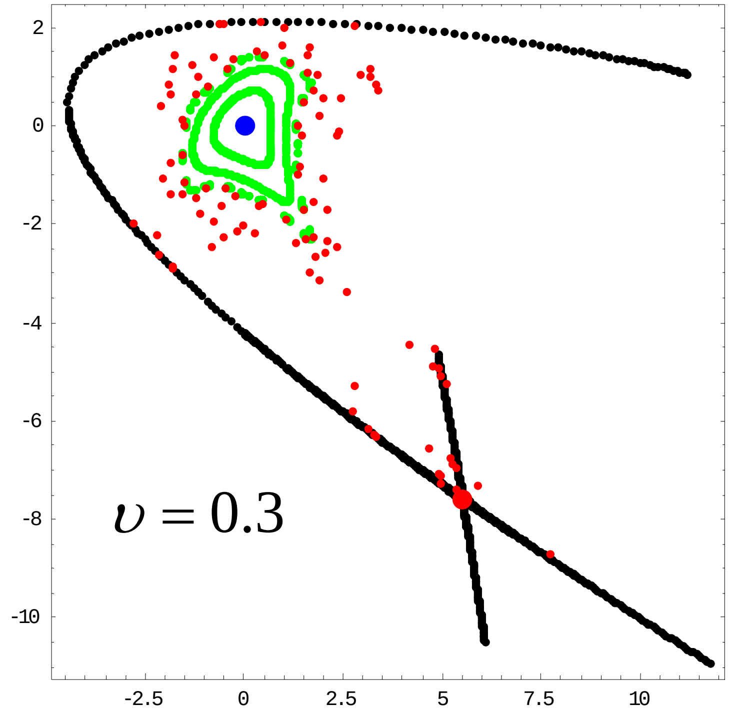
Often the dynamic aperture is much smaller than the fixed point indicates !

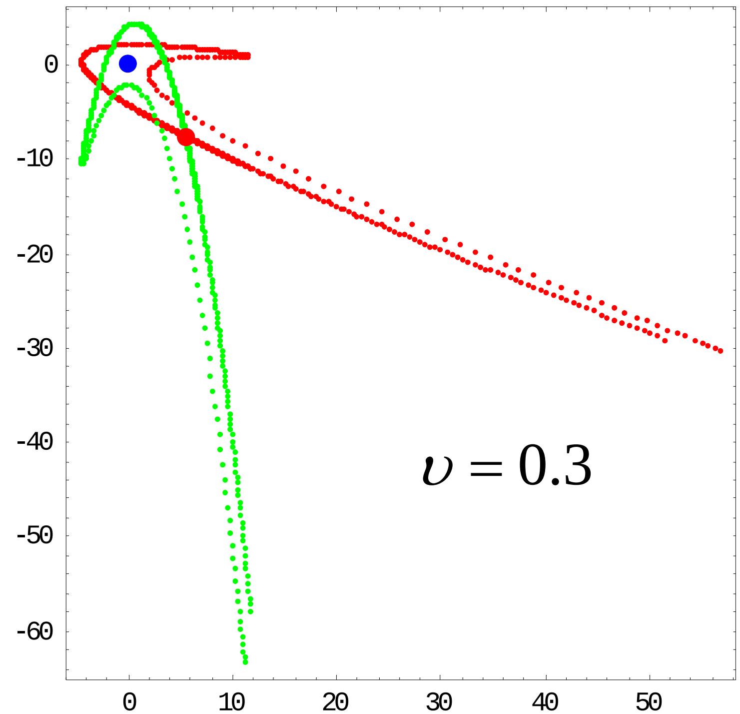
When many sextupoles are used: $$\xi_{0x} + \frac{N}{4\pi}\beta_x\eta_x k_2 l \approx 0$$

The sum of all k_2^2 is then reduced to about $\sum (k_2 l\beta)^2 \approx N(k_2 l\beta)^2 \approx \frac{1}{N}\left(\frac{4\pi}{\eta}\xi_0\right)^2$

The dynamic aperture is therefore greatly increased when distributed sextupoles are used.

Sextupole Extraction



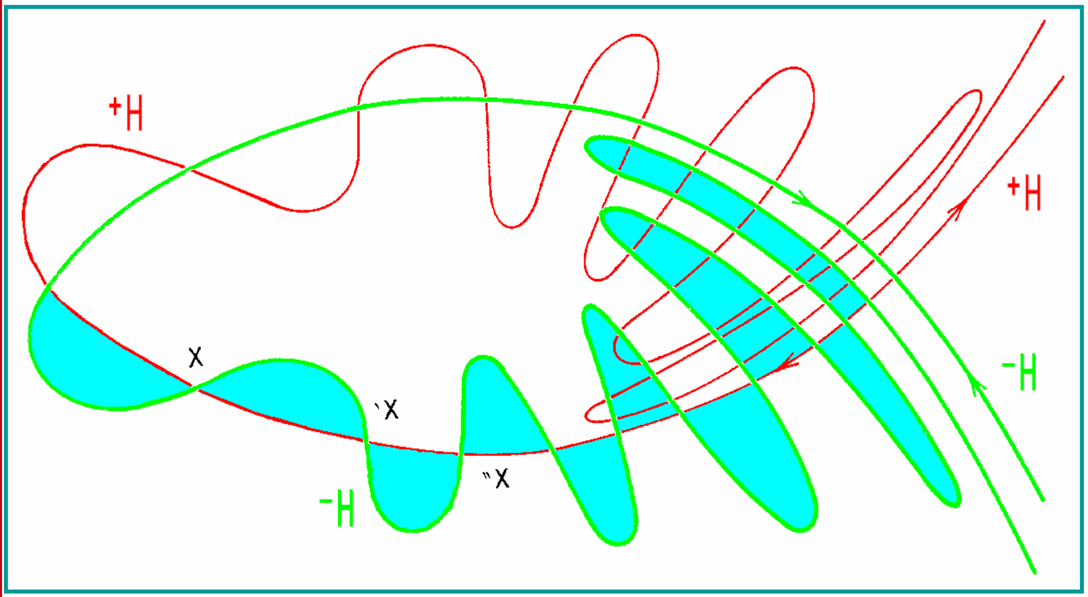


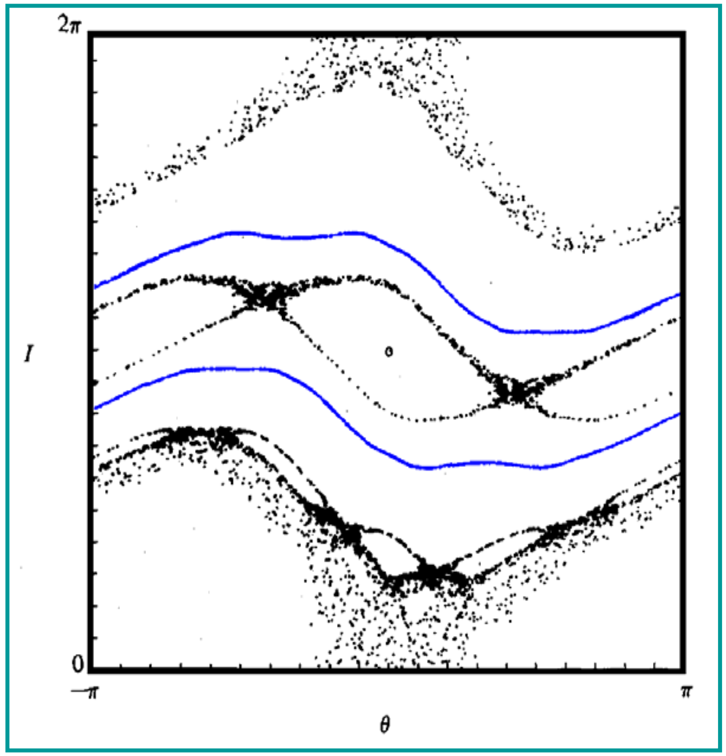
Due to the narrow region of unstable trajectories, sextupoles are used for slow particle extraction at a tune of 1/3.

The intersection of stable and unstable manifolds is a certain indication of chaos.

Homoclinic Points

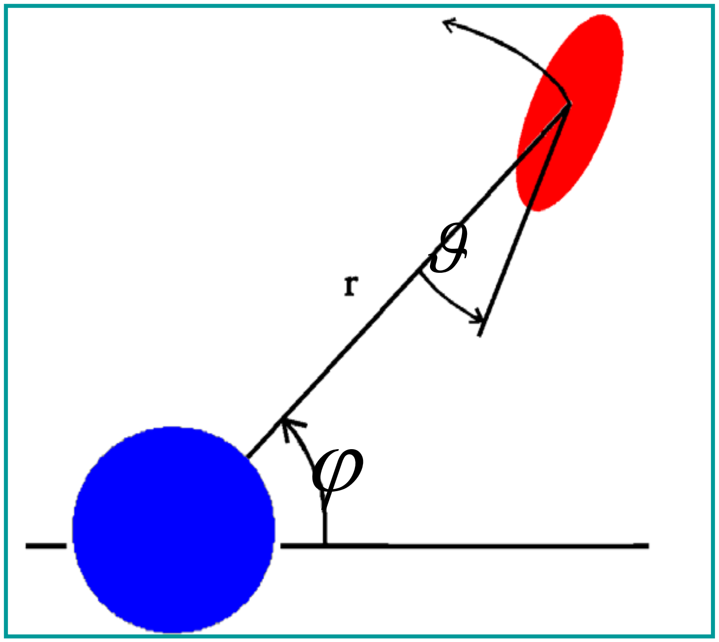
- At instable fixed points, there is a stable and an instabile invariant curve.
- Intersections of these curves (homoclinic points) lead to chaos.

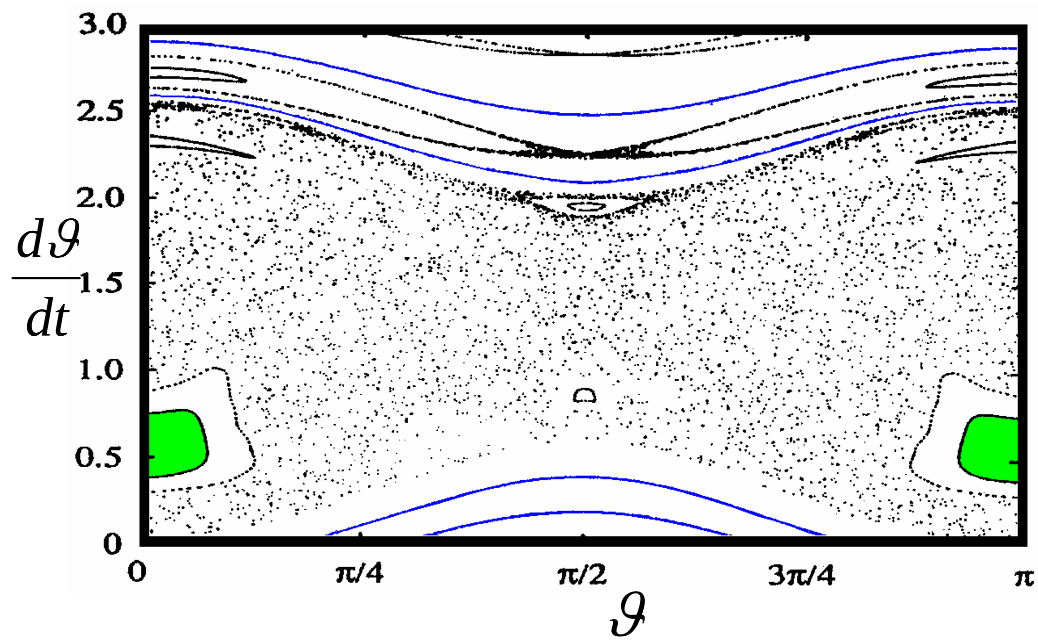




Hyperion: rotation around the vertical

$$\frac{d^2(\vartheta+\varphi(t))}{dt^2} = -\alpha\left(\frac{a}{r(t)}\right)^3 \sin 2\vartheta$$





- On the path from Rotation to Libration around the Spin-Orbit-Coupling is a strong chaotic region.

Tilt of the earth

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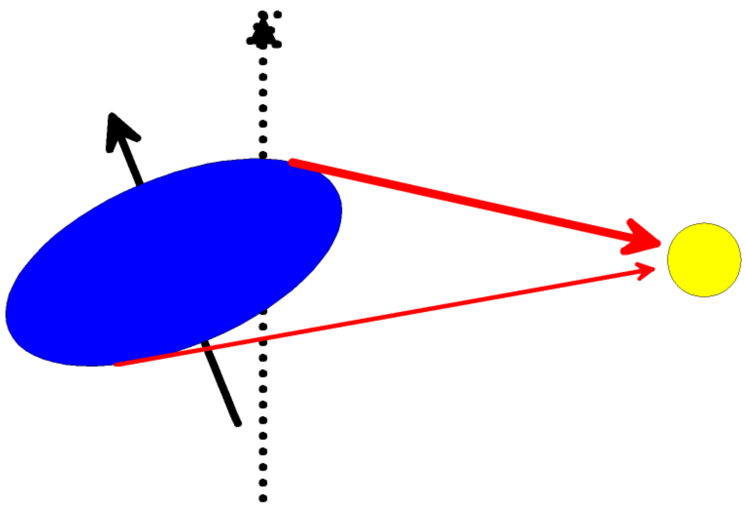
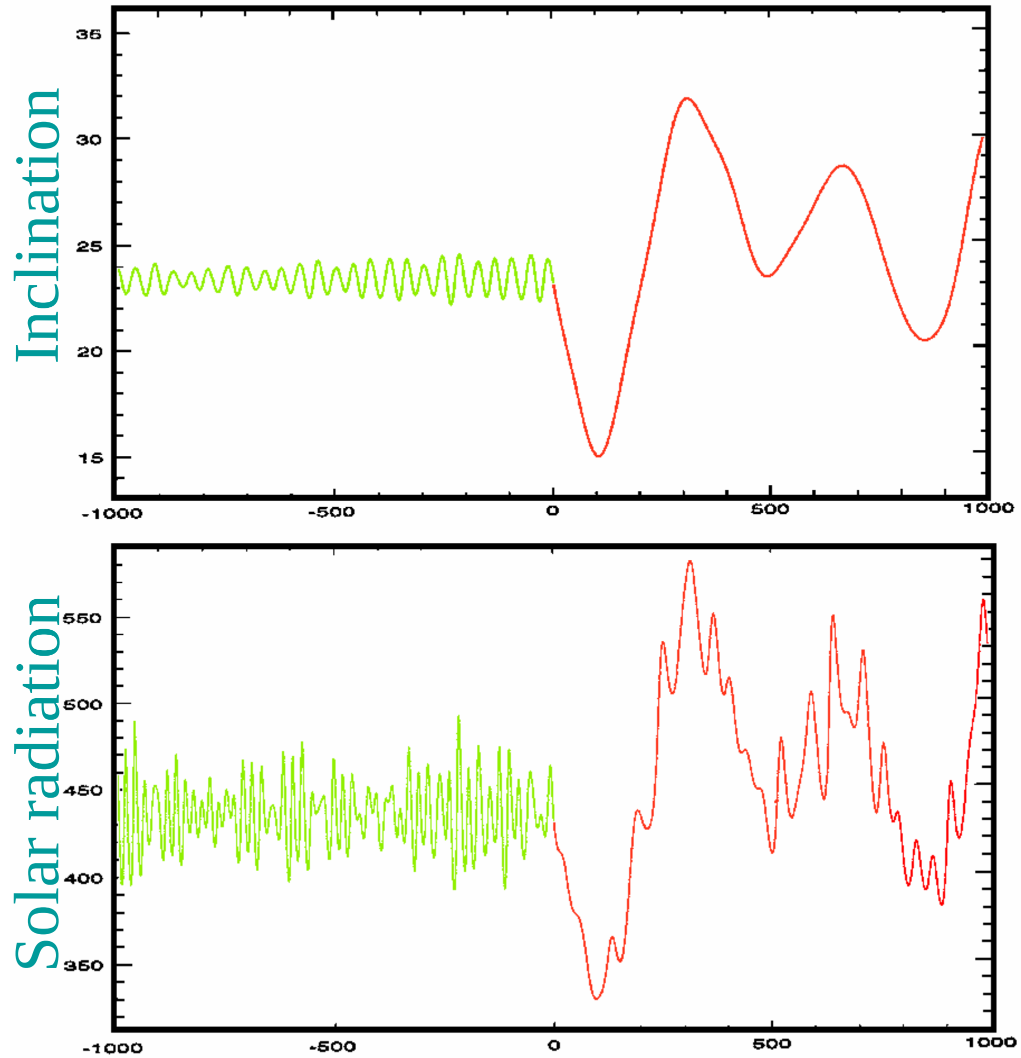


- Tidal forces from moon and sun cause a stabilization of the rotation axis.



-1Ma with and without moon 1Ma

Perturbations

$$\begin{pmatrix} x' \\ a' \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ -K & 0 \end{pmatrix} \begin{pmatrix} x \\ a \end{pmatrix} + \begin{pmatrix} 0 \\ \Delta f \end{pmatrix}$$

$$\begin{pmatrix} x \\ a \end{pmatrix} = \sqrt{2J} \begin{pmatrix} \sqrt{\beta} & 0 \\ -\frac{\alpha}{\sqrt{\beta}} & \frac{1}{\sqrt{\beta}} \end{pmatrix} \begin{pmatrix} \sin(\psi + \phi_0) \\ \cos(\psi + \phi_0) \end{pmatrix} = \sqrt{2J}\, \underline{\beta}\, \vec{S}$$

This would be a solution with constant J and ϕ when $\Delta f = 0$.

Variation of constants:

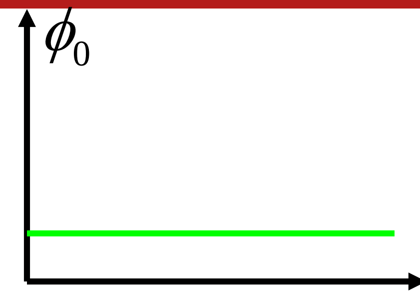
$$\frac{J'}{\sqrt{2J}} \underline{\beta}\, \vec{S} + \sqrt{2J}\, \phi_0' \begin{pmatrix} 0 & \sqrt{\beta} \\ -\frac{1}{\sqrt{\beta}} & -\frac{\alpha}{\sqrt{\beta}} \end{pmatrix} \vec{S} = \begin{pmatrix} 0 \\ \Delta f \end{pmatrix}$$

$$\frac{J'}{\sqrt{2J}} \vec{S} + \sqrt{2J}\, \phi_0' \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \vec{S} = \underline{\beta}^{-1} \begin{pmatrix} 0 \\ \Delta f \end{pmatrix} \quad \text{with} \quad \underline{\beta}^{-1} = \begin{pmatrix} \frac{1}{\sqrt{\beta}} & 0 \\ \frac{\alpha}{\sqrt{\beta}} & \sqrt{\beta} \end{pmatrix}$$

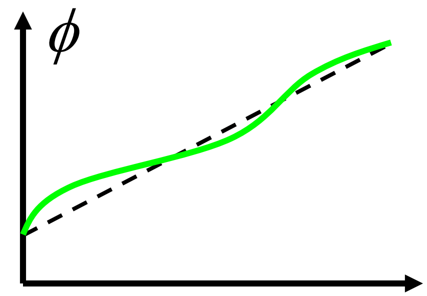
$$\frac{J'}{\sqrt{2J}} = \cos(\psi + \phi_0)\sqrt{\beta}\Delta f \quad , \quad \sqrt{2J}\, \phi_0' = -\sin(\psi + \phi_0)\sqrt{\beta}\Delta f$$

Simplification of linear motion

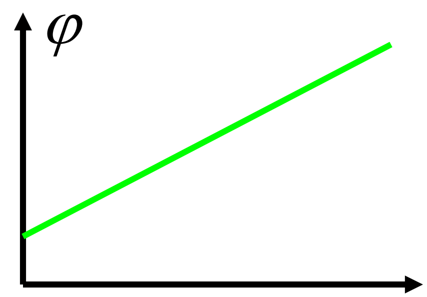
$$\begin{pmatrix} x \\ a \end{pmatrix} = \sqrt{2J} \begin{pmatrix} \sqrt{\beta} & 0 \\ -\frac{\alpha}{\sqrt{\beta}} & \frac{1}{\sqrt{\beta}} \end{pmatrix} \begin{pmatrix} \sin(\psi + \phi_0) \\ \cos(\psi + \phi_0) \end{pmatrix} \Rightarrow \begin{matrix} J' = 0 \\ \phi_0' = 0 \end{matrix}$$



$$\begin{pmatrix} x \\ a \end{pmatrix} = \sqrt{2J} \begin{pmatrix} \sqrt{\beta} & 0 \\ -\frac{\alpha}{\sqrt{\beta}} & \frac{1}{\sqrt{\beta}} \end{pmatrix} \begin{pmatrix} \sin\phi \\ \cos\phi \end{pmatrix} \Rightarrow \begin{matrix} J' = 0 \\ \phi' = \frac{1}{\beta} \end{matrix}$$



$$\begin{pmatrix} x \\ a \end{pmatrix} = \sqrt{2J} \begin{pmatrix} \sqrt{\beta} & 0 \\ -\frac{\alpha}{\sqrt{\beta}} & \frac{1}{\sqrt{\beta}} \end{pmatrix} \begin{pmatrix} \sin(\psi - \mu\frac{s}{L} + \varphi) \\ \cos(\psi - \mu\frac{s}{L} + \varphi) \end{pmatrix} \Rightarrow \begin{matrix} J' = 0 \\ \varphi' = \mu\frac{1}{L} \end{matrix}$$



$$\tilde{\psi} = \psi - \mu\frac{s}{L} \Rightarrow \tilde{\psi}(s+L) = \tilde{\psi}(s)$$

Corresponds to Floquet's Theorem

Quasi-periodic Perturbation

$$J' = \cos(\psi+\phi)\sqrt{2J\beta}\Delta f \quad , \quad \phi' = -\sin(\psi+\phi)\sqrt{\tfrac{\beta}{2J}}\Delta f$$

$$J' = \cos(\tilde{\psi}+\varphi)\sqrt{2J\beta}\Delta f \quad , \quad \varphi' = \mu\tfrac{1}{L} - \sin(\tilde{\psi}+\varphi)\sqrt{\tfrac{\beta}{2J}}\Delta f$$

New independent variable $\vartheta = 2\pi \frac{s}{L}$

$$\tfrac{d}{d\vartheta} J = \cos(\tilde{\psi}+\varphi)\sqrt{2J\beta}\Delta f \, \tfrac{L}{2\pi} \quad , \quad \tfrac{d}{d\vartheta}\varphi = \upsilon - \sin(\tilde{\psi}+\varphi)\sqrt{\tfrac{\beta}{2J}}\Delta f \, \tfrac{L}{2\pi}$$

$$\Delta f(x) = \Delta f(\sqrt{2J\beta}\sin(\tilde{\psi}+\varphi))$$

The perturbations are 2π periodic in ϑ and in φ

φ is approximately $\varphi \approx \upsilon \cdot \vartheta$

For irrational ν, the perturbations are quasi-periodic.

Tune Shift with Amplitude

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$$\tfrac{d}{d\vartheta}J = \cos(\tilde{\psi}+\varphi)\sqrt{2J\beta}\,\Delta f\,\tfrac{L}{2\pi} \quad , \quad \tfrac{d}{d\vartheta}\varphi = \upsilon - \sin(\tilde{\psi}+\varphi)\sqrt{\tfrac{\beta}{2J}}\,\Delta f\,\tfrac{L}{2\pi}$$

$$\tfrac{d}{d\vartheta}\varphi = \partial_J H \quad , \quad \tfrac{d}{d\vartheta}J = -\partial_\phi H \quad , \quad H(\varphi, J, \vartheta) = \upsilon \cdot J - \tfrac{L}{2\pi}\int_0^x \Delta f(\hat{x}, s)\,d\hat{x}$$

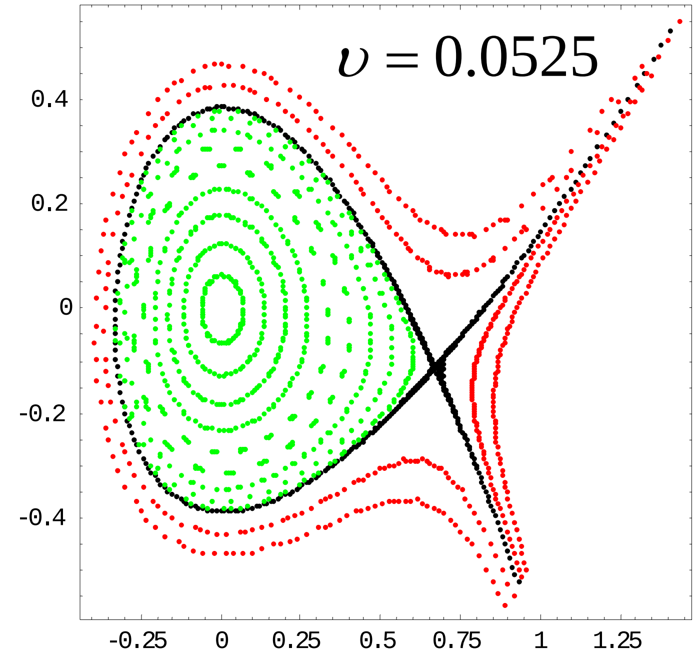
The motion remains Hamiltonian in the perturbed coordinates !

If there is a part in $\partial_J H$ that does not depend on φ, s $\Rightarrow$ Tune shift

The effect of other terms tends to average out.

$$\varphi(\vartheta) - \varphi_0 \approx \vartheta \cdot \partial_J \langle H \rangle_{\varphi,\vartheta}(J)$$

$$\upsilon(J) = \upsilon + \partial_J \langle \Delta H \rangle_{\varphi,\vartheta}(J)$$



Tune Shift Examples

$$H(\varphi,J)=\upsilon\cdot J-\tfrac{L}{2\pi}\int_0^x \Delta f(\hat{x},s)\,d\hat{x} \quad , \quad \Delta\upsilon(J)=\partial_J\langle\Delta H\rangle_{\varphi,\vartheta}$$

Quadrupole: $\Delta f = -\Delta k\, x$

$$\Delta H = \tfrac{L}{2\pi}\Delta k\,\tfrac{1}{2}x^2 = \tfrac{L}{2\pi}\Delta k\, J\beta\sin^2(\tilde{\psi}+\varphi)$$

$$\langle\Delta H\rangle_{\varphi,\vartheta} = \tfrac{1}{2\pi}\int_0^{2\pi}\Delta k\beta\, d\vartheta\, L\frac{J}{4\pi} = \int_0^L \Delta k\beta\, ds\frac{J}{4\pi} \Rightarrow \Delta\upsilon = \tfrac{1}{4\pi}\oint\Delta k\beta\, ds$$

Sextupole: $\Delta f = -k_2\tfrac{1}{2}x^2$

$$\Delta H = \tfrac{L}{2\pi}k_2\tfrac{1}{3!}x^3 = \tfrac{L}{2\pi}k_2\tfrac{1}{3!}\sqrt{2J\beta}^{\,3}\sin^3(\tilde{\psi}+\varphi)$$

$$\langle\Delta H\rangle_{\varphi,\vartheta} = 0 \quad \Rightarrow \quad \Delta\upsilon = 0$$

Octupole: $\Delta f = -k_3\tfrac{1}{3!}x^3$

$$\Delta H = \tfrac{L}{2\pi}k_3\tfrac{1}{4!}x^4 = \tfrac{L}{2\pi}k_3\tfrac{1}{3!}(J\beta)^2\sin^4(\tilde{\psi}+\varphi)$$

$$\langle\Delta H\rangle_{\varphi,\vartheta} = \tfrac{J^2}{3!2\pi}\oint k_3\beta^2 ds\left\langle\tfrac{1}{2^4}(e^{i\varphi}-e^{-i\varphi})^4\right\rangle_{\varphi} \Rightarrow \Delta\upsilon = J\,\tfrac{1}{16\pi}\oint k_3\beta^2 ds$$

Nonlinear Resonances

$$\frac{d}{d\vartheta}J = \cos(\tilde{\psi}+\varphi)\sqrt{2J\beta}\Delta f\,\frac{L}{2\pi} \quad , \quad \frac{d}{d\vartheta}\varphi = \upsilon - \sin(\tilde{\psi}+\varphi)\sqrt{\frac{\beta}{2J}}\Delta f\,\frac{L}{2\pi}$$

$$\frac{d}{d\vartheta}\varphi = \partial_J H \quad , \quad \frac{d}{d\vartheta}J = -\partial_\phi H \quad , \quad H(\varphi, J, \vartheta) = \upsilon\cdot J - \frac{L}{2\pi}\int_0^x \Delta f(\hat{x}, s)\,d\hat{x}$$

The effect of the perturbation is especially strong when

$$\cos(\tilde{\psi}+\varphi)\sqrt{\beta}\Delta f \quad \text{or} \quad \sin(\tilde{\psi}+\varphi)\sqrt{\beta}\Delta f$$

has contributions that hardly change, i.e. the change of

$\sqrt{\beta(\vartheta)}\Delta f(x(\vartheta),\vartheta)$ is in resonance with the rotation angle $\varphi(\vartheta)$.

Periodicity allows Fourier expansion:

$$H(\varphi, J, \vartheta) = \sum_{n,m=-\infty}^{\infty} \widehat{H}_{nm}(J)\,e^{i[n\vartheta+m\varphi]} = \sum_{n,m=-\infty}^{\infty} H_{nm}(J)\cos(n\vartheta + m\varphi + \Psi_{nm}(J))$$

$$H_{00}(J) = \langle H(\varphi, J, s)\rangle_{\varphi,s} \Rightarrow$$ Tune shift

The Single Resonance Model

$$\frac{d}{d\vartheta}J = \sum_{n,m=-\infty}^{\infty} mH_{nm}(J)\sin(n\vartheta + m\varphi + \Psi_{nm}(J))$$

$$\frac{d}{d\vartheta}\varphi = \upsilon + \partial_J \sum_{n,m=-\infty}^{\infty} H_{nm}(J)\cos(n\vartheta + m\varphi + \Psi_{nm}(J))$$

Strong deviation from: $J = J_0\ ,\quad \varphi = \upsilon\vartheta + \varphi_0$

Occur when there is coherence between the perturbation and the phase space rotation: $n + m\frac{d}{ds}\varphi \approx 0$

Resonance condition: tune is rational $\boxed{n + m\,\upsilon = 0}$

On resonance the integral would increases indefinitely !

Neglecting all but the most important term

$$H(\varphi, J, \vartheta) \approx \upsilon J + H_{00}(J) + H_{nm}(J)\cos(n\vartheta + m\varphi + \Psi_{nm}(J))$$

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Fixed points

$$\frac{d}{d\vartheta} J = mH_{nm}(J)\sin(n\vartheta + m\varphi + \Psi_{nm}(J))$$

$$\frac{d}{d\vartheta} \varphi = \upsilon + \Delta\upsilon(J) + \partial_J [H_{nm}(J)\cos(n\vartheta + m\varphi + \Psi_{nm}(J))]$$

$$\Phi = \tfrac{1}{m}[n\vartheta + m\varphi + \Psi_{nm}(J)] \ , \quad \delta = \upsilon + \tfrac{n}{m}$$

$$\frac{d}{d\vartheta} J = mH_{nm}(J)\sin(m\Phi) \ , \quad \frac{d}{d\vartheta}\Phi = \delta + \Delta\upsilon(J) + H'_{nm}(J)\cos(m\Phi)$$

$$H(\varphi, J, \vartheta) \approx \delta J + H_{00}(J) + H_{nm}(J)\cos(m\Phi)$$

Fixed points: $\frac{d}{d\vartheta} J = mH_{nm}(J_f)\sin(m\Phi_f) = 0 \quad \Rightarrow \quad \Phi_f = \frac{k}{m}\pi$

If $\delta + \Delta\upsilon(J_f) \pm H'_{nm}(J_f) = 0$ has a solution.

$$\frac{d}{d\vartheta}\Delta J = \pm m^2 H_{nm}(J_f)\Delta\Phi \ , \quad \frac{d}{d\vartheta}\Delta\Phi = [\Delta\upsilon'(J_f) \pm H''_{nm}(J_f)]\Delta J$$

Stable fixed point for: $H_{nm}(J_f)[H''_{nm}(J_f) \pm \Delta\upsilon'(J_f)] < 0$

Third Integer Resonances

Sextupole: $\Delta f = -k_2 \frac{1}{2} x^2$

$$\Delta H = \frac{L}{2\pi} k_2 \frac{1}{3!} x^3 = \frac{L}{2\pi} k_2 \frac{1}{3!} \sqrt{2J\beta}^3 \sin^3(\tilde{\psi} + \varphi)$$

$$= \frac{L}{2\pi} k_2 \frac{1}{3!4} \sqrt{2J\beta}^3 [\sin(3[\tilde{\psi} + \varphi]) + 3\sin(\tilde{\psi} + \varphi)]$$

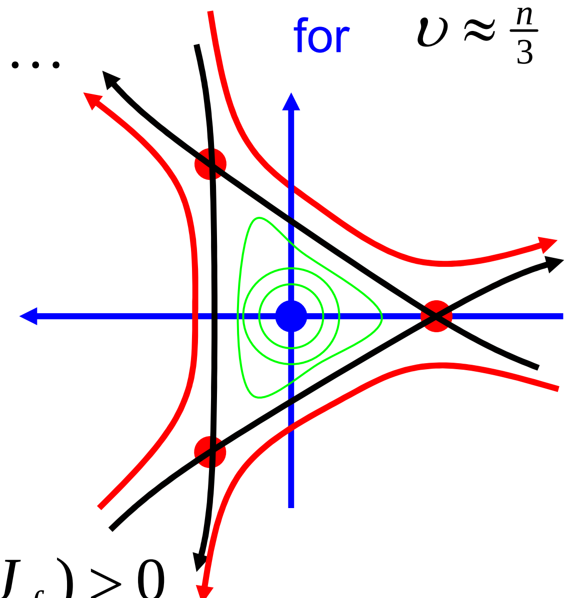
Simplification: one sextupole $k_2(\vartheta) = k_2\delta(\vartheta) = k_2 \frac{1}{2\pi} \sum_{n=-\infty}^{\infty} \cos(n\vartheta)$

$$\Delta H = \frac{L}{2\pi} k_2 \frac{1}{3!4} \sqrt{2J\beta}^3 \frac{1}{2\pi} \cos(-n\vartheta + 3\varphi + \tilde{\psi} - \tfrac{\pi}{2}) + \dots$$

for $\upsilon \approx \frac{n}{3}$

$$\Delta H \approx A_2 \sqrt{J}^3 \cos(3\Phi)$$

$$\left.\begin{array}{l} \Phi_f = 0, \frac{1}{3}\pi, \frac{2}{3}\pi, \dots \\ \delta \pm A_2 \frac{3}{2}\sqrt{J} = 0 \end{array}\right\} \begin{array}{l} \Phi_f = \frac{1}{3}\pi, \pi, \frac{5}{3}\pi \\ \text{for } \delta > 0 \end{array}$$



All these fixed points are instable since $H_{nm}(J_f)\, H''_{nm}(J_f) > 0$

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Fourth Integer Resonances

Octupole: $\boxed{\Delta f = -k_3 \frac{1}{3!} x^3}$, $\Delta H = \frac{L}{2\pi} k_3 \frac{1}{4!} x^4 = \frac{L}{2\pi} k_3 \frac{1}{3!} J^2 \beta^2 \sin^4(\tilde{\psi} + \varphi)$

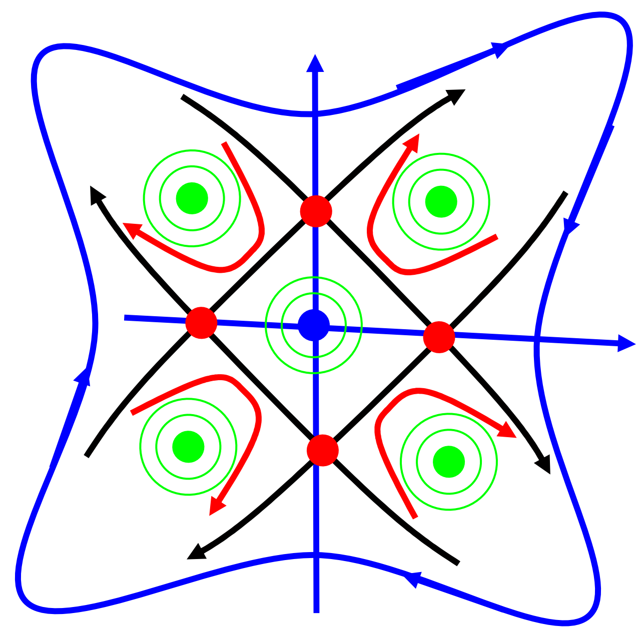
$$= \frac{L}{2\pi} k_3 \frac{1}{3!8} J^2 \beta^2 [\cos(4[\tilde{\psi} + \varphi]) - 4\cos(\tilde{\psi} + \varphi) + 3]$$

Simplification: one octupole $k_3(\vartheta) = k_3 \delta(\vartheta) = k_3 \frac{1}{2\pi} \sum_{n=-\infty}^{\infty} \cos(n\vartheta)$

$\Delta H \approx A_3 J^2 [3 + \cos(4\Phi)]$ for $\upsilon \approx \frac{n}{4}$

$\Phi_f = 0, \frac{1}{4}\pi, \frac{2}{4}\pi, \ldots$ Either 8 fixed points: $\delta < 0$

$\delta + A_3 2J(3 \pm 1) = 0$ or none for: $\delta > 0$



$$H_{nm}(J_f)[H''_{nm}(J_f) \pm \Delta\upsilon'(J_f)] < 0$$

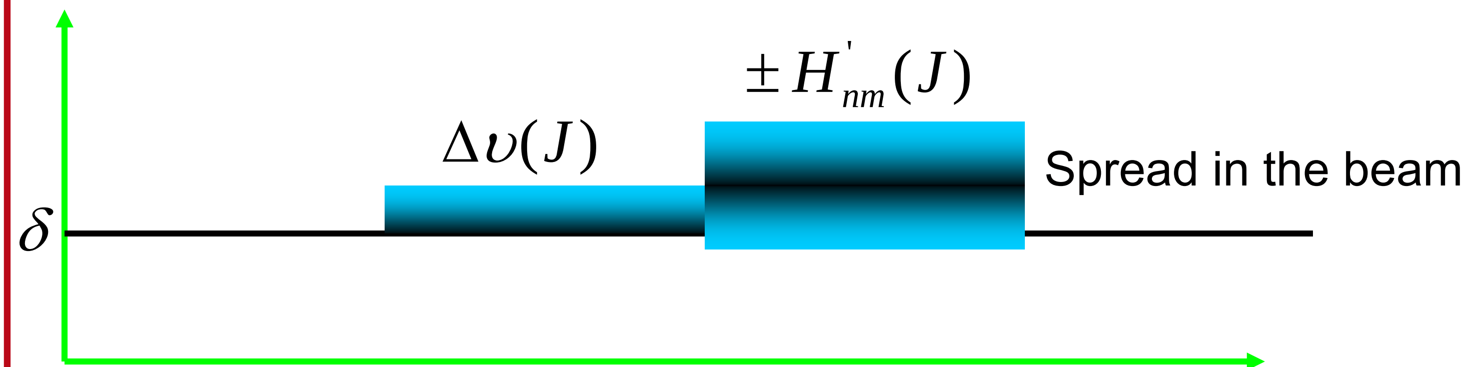
Stability for $(2A_3 J)^2 [1 \pm 3] < 0$,
i.e. for the 4 outer fixed points.

Resonance Width (Strength)

Fixed points: $\frac{d}{d\vartheta} J = mH_{nm}(J_f)\sin(m\Phi_f) = 0 \quad \Rightarrow \quad \Phi_f = \frac{k}{m}\pi$

If $\delta + \Delta\upsilon(J_f) \pm H'_{nm}(J_f) = 0$ has a solution.

δ has to avoid the region $\delta + \Delta\upsilon(J) \pm H'_{nm}(J) = 0$ for all particles.



Assuming that the tune shift and perturbation are monotonous in J:

This tune region has the width $\Delta_{nm} = 2|H'_{nm}(J_{\max})|$ for strong resonances.

Δ_{nm} Is called Resonance Width, Resonance Strength, or Stop-Band Width

Coupling Resonances

$$\frac{d}{d\vartheta} J_x = \cos(\tilde{\psi}_x + \varphi_x)\sqrt{2J_x\beta_x}\Delta f_x \frac{L}{2\pi} \quad , \quad \frac{d}{d\vartheta}\varphi_x = \upsilon_x - \sin(\tilde{\psi}_x + \varphi_x)\sqrt{\frac{\beta_x}{2J_x}}\Delta f_x \frac{L}{2\pi}$$

$$\frac{d}{d\vartheta} J_y = \cos(\tilde{\psi}_y + \varphi_y)\sqrt{2J_y\beta_y}\Delta f_y \frac{L}{2\pi} \quad , \quad \frac{d}{d\vartheta}\varphi_y = \upsilon_y - \sin(\tilde{\psi}_y + \varphi_y)\sqrt{\frac{\beta_y}{2J_y}}\Delta f_y \frac{L}{2\pi}$$

$$\frac{d}{d\vartheta}\vec{\varphi} = \vec{\partial}_J H \quad , \quad \frac{d}{d\vartheta}\vec{J} = -\vec{\partial}_\varphi H \quad , \quad H(\vec{\varphi},\vec{J},\vartheta) = \vec{\upsilon}\cdot\vec{J} - \frac{L}{2\pi}\int_0^{\vec{x}} \Delta\vec{f}(\hat{\vec{x}},s)d\hat{\vec{x}}$$

The integral form can be chosen since it is path independent. This is due to the Hamiltonian nature of the force: $\Delta f_{x,y}(x,y,s) = -\partial_{x,y}\Delta H(x,y,s)$

Single Resonance model for two dimensions means retaining only the amplitude dependent tune shift and one term in the two dimensional Fourier expansion:

$$H(\vec{\varphi},\vec{J},\vartheta) = \vec{\upsilon}\cdot\vec{J} + H_{00}(\vec{J}) + H_{n\vec{m}}(\vec{J})\cos(n\vartheta + m_x\varphi_x + m_y\varphi_y + \Psi_{n\vec{m}}(\vec{J}))$$

For $n + m_x\upsilon_x + m_y\upsilon_y \approx 0$

$$m_x\varphi_x + m_y\varphi_y = \vec{m}\cdot\vec{\varphi}$$

Sum and Difference Resonances

$n+m_x\upsilon_x+m_y\upsilon_y\approx 0$ means that oscillations in y can drive oscillations in x in

$$x''=-K\,x+\Delta f_x(x,y,s)$$

The resonance term in the Hamiltonian then changes only slowly:

$$H(\vec{\varphi},\vec{J},\vartheta)=\vec{\upsilon}\cdot\vec{J}+H_{00}(\vec{J})+H_{n\vec{m}}(\vec{J})\cos(n\vartheta+m_x\varphi_x+m_y\varphi_y+\Psi_{n\vec{m}}(\vec{J}))$$

$$\tfrac{d}{d\vartheta}\vec{\varphi}=\vec{\partial}_J H \quad , \quad \tfrac{d}{d\vartheta}\vec{J}=-\vec{\partial}_\varphi H$$

$$J=\vec{m}\cdot\vec{J}$$

$$J_\perp=m_xJ_x-m_yJ_y=\vec{m}\times\vec{J} \quad\Rightarrow\quad \tfrac{d}{d\vartheta}J_\perp=0$$

Difference resonances lead to stable motion since:

$$n+|m_x|\upsilon_x-|m_y|\upsilon_y\approx 0\Rightarrow|m_x|J_x+|m_y|J_y=const.$$

Sum resonances lead to unstable motion since:

$$n+|m_x|\upsilon_x+|m_y|\upsilon_y\approx 0\Rightarrow|m_x|J_x-|m_y|J_y=const.$$

Resonances Diagram

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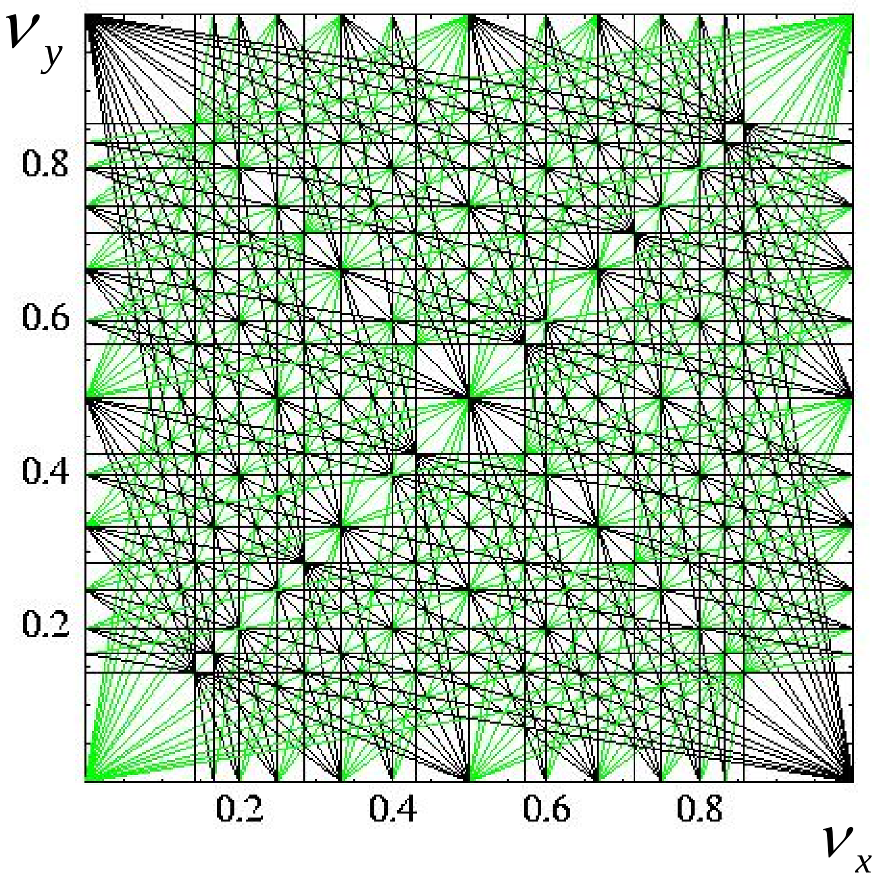
$n + m_x \upsilon_x + m_y \upsilon_y \approx 0$ means that oscillations in y can drive oscillations in x in

$$x'' = -K\,x + \Delta f_x(x, y, s)$$



All these resonances have to be avoided by their respective resonance width.

The position of an accelerator in the tune plane s called its Working Point.