

Expectation values for p_x

Review of one dimension:

$\Psi_{n+1}(x_{n+1} - \frac{\Delta}{2}) = \Psi_n(x_n + \frac{\Delta}{2})$

$$\Psi_n = A_n e^{i(k_n x_n - \omega t)}$$

$$\Psi_{n+1} = A_{n+1} e^{i(k_{n+1} x_{n+1} - \omega t)}$$

$$A_{n+1} = A_n e^{ik_n(x_n + \frac{\Delta}{2})} e^{-ik_{n+1}(x_{n+1} - \frac{\Delta}{2})}$$

$$\Psi_{n+1} = A_n e^{i[k_n(x_n + \frac{\Delta}{2}) - k_{n+1}(x_{n+1} - \frac{\Delta}{2}) + k_{n+1}x_{n+1}]} e^{-i\omega t} = \Psi_n e^{i[k_n \frac{\Delta}{2} + k_{n-1} \frac{\Delta}{2}]}$$

$$= \Psi_n e^{i \frac{k_{n+1} + k_n}{2} \Delta} \approx \Psi_n + ik_n \Psi_n \Delta \quad \rightarrow \quad \underline{\underline{\frac{\partial}{\partial x} \Psi(x, t) = ik \Psi}}$$

Conclusion: whenever $k \Psi$ needs to be computed, one can use $-i \frac{\partial}{\partial x} \Psi$

$$\langle p_x \rangle = \sum_{\text{all } j} \hbar k_j |\Psi_j|^2 dx = \sum_{\text{all } j} \Psi_j^* \left(-i \hbar \frac{\partial}{\partial x} \Psi(x_j) \right) dx = \int_{-\infty}^{\infty} \Psi^* \left(-i \hbar \frac{\partial}{\partial x} \right) \Psi dx$$



Spread of measurements

$$\Delta x^2 = \int_{-\infty}^{\infty} \Psi^*(x,t) (x - \langle x \rangle)^2 \Psi(x,t) dx = \langle x^2 - 2x\langle x \rangle + \langle x \rangle^2 \rangle = \langle x^2 \rangle - \langle x \rangle^2$$

$$\Delta p_x^2 = \int_{-\infty}^{\infty} \Psi^*(x,t) (\hat{p}_x - \langle p_x \rangle)^2 \Psi(x,t) dx = \langle p_x^2 \rangle - \langle p_x \rangle^2, \quad \hat{p}_x = -i\hbar \frac{\partial}{\partial x}$$

$$\Delta E^2 = \int_{-\infty}^{\infty} \Psi^*(x,t) (\hat{E} - \langle E \rangle)^2 \Psi(x,t) dx = \langle E^2 \rangle - \langle E \rangle^2, \quad \hat{E} = V(x) - \frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2}$$

For an arbitrary physical quantity **A** that depends on x and p_x , a corresponding operator can be formed. The spread of measurements of **A** is then given by:

$$\Delta A = \int_{-\infty}^{\infty} \Psi^*(x,t) (\hat{A} - \langle A \rangle)^2 \Psi(x,t) dx = \langle A^2 \rangle - \langle A \rangle^2$$

The spread of the measurements is 0 when the wave function is an eigenfunction of the operator.

$$\int_{-\infty}^{\infty} \Psi_a^*(x,t) (\hat{A} - \langle A \rangle)^2 \Psi_a(x,t) dx = 0 \quad \text{for} \quad \hat{A} \Psi_a(x,t) = a \Psi_a(x,t)$$

since $\langle A \rangle = a$ and $\langle A^2 \rangle = a^2$

