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P443 Prelim II
April 11, 2008

1. Spin 1/2

Consider a spin 1/2 particle with magnetic moment $\vec{\mu} = \frac{e}{m}\vec{S}$ in a magnetic field $\vec{B} = B_0\hat{z}$. The hamiltonian

$$H = -\vec{\mu} \cdot \vec{B} = \frac{e}{m}\vec{S} \cdot \vec{B}$$

where $\vec{S} = \frac{\hbar}{2}\vec{\sigma}$. For convenience we define the constant $\gamma = \frac{e\hbar}{2m}$.

- (a) Find the eigenvectors (wave functions) and the eigenvalues (energy levels) of the hamiltonian.

[The hamiltonian is

$$H = \frac{e}{m}S_z B_0 = \gamma \begin{pmatrix} B_0 & 0 \\ 0 & -B_0 \end{pmatrix}.$$

The eigenvectors are $\chi_+ = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$, $\chi_- = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$ with eigenvalues $\pm\gamma B_0$.]

Now suppose we add a magnetic field in the x-direction such that $\vec{B} = B_0\hat{z} + B_x\hat{x}$ and $B_x \ll B_0$.

- (b) Use perturbation theory to calculate the first order shift in the energy of each of the eigenstates due to $H' = -\vec{\mu} \cdot B_x\hat{x}$.

[The perturbation is

$$H' = \frac{e}{m}S_x B_x = \gamma \begin{pmatrix} 0 & B_x \\ B_x & 0 \end{pmatrix}$$

The first order shifts are

$$E_+^1 = \langle \chi_+ | H' | \chi_+ \rangle = (1 \ 0) \gamma \begin{pmatrix} 0 & B_x \\ B_x & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = 0$$
$$E_-^1 = \langle \chi_- | H' | \chi_- \rangle = (0 \ 1) \gamma \begin{pmatrix} 0 & B_x \\ B_x & 0 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = 0$$

- (c) Use perturbation theory to compute the second order shift in the energy of each of the eigenstates.

[The second order energy shift

$$E_+^2 = \frac{|\langle \chi_- | H' | \chi_+ \rangle|^2}{E_+^0 - E_-^0} = \frac{1}{2\gamma B_0} |(0 \ 1) \gamma \begin{pmatrix} 0 & B_x \\ B_x & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix}|^2 = \frac{\gamma B_x^2}{2B_0}$$

$$E_-^2 = \frac{|\langle \chi_+ | H' | \chi_- \rangle|^2}{E_-^0 - E_+^0} = \frac{1}{-2\gamma B_0} |(1 \ 0) \gamma \begin{pmatrix} 0 & B_x \\ B_x & 0 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix}|^2 = -\frac{\gamma B_x^2}{2B_0}]$$

- (d) Use perturbation theory to compute the first order shift in the wave function of each of the eigenstates

[The first order shift in the wave functions is

$$\chi_+^2 = \frac{\langle \chi_- | H' | \chi_+ \rangle}{E_+^0 - E_-^0} \chi_- = \frac{1}{2\gamma B_0} (0 \ 1) \gamma \begin{pmatrix} 0 & B_x \\ B_x & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} \chi_- = \frac{B_x}{2B_0} \chi_-$$

$$\chi_-^2 = \frac{\langle \chi_+ | H' | \chi_- \rangle}{E_-^0 - E_+^0} \chi_+ = -\frac{1}{2\gamma B_0} (1 \ 0) \gamma \begin{pmatrix} 0 & B_x \\ B_x & 0 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} \chi_+ = -\frac{B_x}{2B_0} \chi_+]$$

- (e) Compute the exact energy levels of the system, expand to second order and compare with the result from perturbation theory.

[The exact hamiltonian is $H = \gamma \begin{pmatrix} B_0 & B_x \\ B_x & -B_0 \end{pmatrix}$. The exact energy levels are the eigenvalues of the hamiltonian, namely

$$E_{\pm} = \pm \gamma \sqrt{B_0^2 + B_x^2} \sim \pm B_0 \left(1 + \frac{1}{2} \frac{B_x^2}{B_0^2} \right).$$

The exact and perturbative solutions are in agreement to second order.]

- (f) Compute the exact wave functions, expand to second order and compare with the result from perturbation theory. (You do not need to normalize.)

[The exact wave functions are the eigenvectors of the exact hamiltonian.

$$\gamma \begin{pmatrix} B_0 & B_x \\ B_x & -B_0 \end{pmatrix} \begin{pmatrix} 1 \\ a \end{pmatrix} = \gamma \sqrt{B_0^2 + B_x^2} \begin{pmatrix} 1 \\ a \end{pmatrix}$$

$$\rightarrow \begin{pmatrix} B_0 + aB_x \\ B_x - aB_0 \end{pmatrix} = \sqrt{B_0^2 + B_x^2} \begin{pmatrix} 1 \\ a \end{pmatrix}$$

$$\begin{aligned}
\rightarrow a &= \frac{\sqrt{B_0^2 + B_x^2}}{B_x} - \frac{B_0}{B_x} \\
&\sim \frac{1}{B_x} \left(B_0 \left(1 + \frac{B_x^2}{2B_0^2} \right) \right) - \frac{B_0}{B_x} \\
&\sim \frac{B_x}{2B_0}
\end{aligned}$$

Then the eigenvector is

$$\begin{pmatrix} 1 \\ \frac{B_x}{2B_0} \end{pmatrix} = \chi_+^0 + \frac{B_x}{2B_0} \chi_-^0$$

in agreement with the perturbation theory result in part (d). Similarly, to get the other eigenvector

$$\begin{aligned}
\gamma \begin{pmatrix} B_0 & B_x \\ B_x & -B_0 \end{pmatrix} \begin{pmatrix} a \\ 1 \end{pmatrix} &= -\gamma \sqrt{B_0^2 + B_x^2} \begin{pmatrix} a \\ 1 \end{pmatrix} \\
\rightarrow \begin{pmatrix} aB_0 + B_x \\ aB_x - B_0 \end{pmatrix} &= -\sqrt{B_0^2 + B_x^2} \begin{pmatrix} a \\ 1 \end{pmatrix} \\
\rightarrow a &= \frac{-\sqrt{B_0^2 + B_x^2}}{B_x} = \frac{B_0}{B_x} \\
&\sim -\frac{1}{B_x} \left(B_0 \left(1 + \frac{B_x^2}{2B_0^2} \right) \right) + \frac{B_0}{B_x} \\
&\sim -\frac{B_x}{2B_0}
\end{aligned}$$

Then the other eigenvector is

$$\begin{pmatrix} -\frac{B_x}{2B_0} \\ 1 \end{pmatrix} = \chi_-^0 - \frac{B_x}{2B_0} \chi_+^0$$

2. Fermi Energy

- (a) Calculate the Fermi energy for electrons in a two-dimensional infinite square well of area A . Give your answer in terms of the density of electrons, (the number of electrons per unit area).

[In a 2-d well, the energy

$$E_F = \frac{\hbar^2 \pi^2}{2m l^2} (n_x^2 + n_y^2)$$

The number of states with energy less than E_F is

$$N_{states} = \frac{1}{4} \pi (n_x^2 + n_y^2) = \frac{1}{4} \pi \frac{E_F 2m l^2}{\hbar^2 \pi^2}$$

There are 2 electrons per state so the total number of electrons is

$$N_{ele} = \frac{1}{2} \pi \frac{E_F 2m l^2}{\hbar^2 \pi^2}$$

The density of electrons (number/area) is

$$\frac{N_{ele}}{l^2} = \sigma = \frac{1}{2} \pi \frac{E_F 2m}{\hbar^2 \pi^2}$$

Then

$$E_F = \frac{\sigma \hbar^2 \pi}{m}]$$

- (b) Calculate the total energy of the electrons

[The total energy is

$$\begin{aligned} E &= \int_0^{E_F} E \frac{dN}{dE} dE = \int_0^{E_F} E \frac{m l^2}{\hbar^2 \pi} dE \\ &= \frac{1}{2} E_F^2 \frac{m l^2}{\hbar^2 \pi} = \frac{1}{2} \left(\frac{\sigma \hbar^2 \pi}{m} \right)^2 \frac{m A}{\hbar^2 \pi} = \frac{1}{2} \frac{(N_{ele}/A)^2 \hbar^2 \pi A}{m} \\ &= \frac{1}{2} \frac{N_{ele}^2 \hbar^2 \pi}{m A} \end{aligned}$$

- (c) Compute the pressure on the walls, $P = \frac{dE_{tot}}{dA}$. The pressure is

$$P = \frac{dE_{tot}}{dA} = -\frac{1}{2} \frac{N_{ele}^2 \hbar^2 \pi}{m A^2} = -\frac{E_{tot}}{A} = -\frac{1}{2} \frac{\sigma^2 \hbar^2 \pi}{m}]$$

3. Three Identical Particles

Imagine a situation in which there are three particles and only three states a , b , and c available to them. If the particles are distinguishable, there are a total of 27 distinct configurations of this system.

What is the total number of allowed distinct configurations for this system

- (a) if the particles are bosons?

[If the particles are bosons there are 10 states:

$$\begin{aligned} &\psi_a(1)\psi_a(2)\psi_a(3), \psi_b(1)\psi_b(2)\psi_b(3), \psi_c(1)\psi_c(2)\psi_c(3), \\ &\psi_a(1)\psi_a(2)\psi_b(3), \psi_a(1)\psi_a(2)\psi_c(3), \psi_b(1)\psi_b(2)\psi_a(3), \\ &\psi_b(1)\psi_b(2)\psi_c(3), \psi_c(1)\psi_c(2)\psi_a(3), \psi_c(1)\psi_c(2)\psi_b(3), \\ &\psi_a(1)\psi_b(2)\psi_c(3)] \end{aligned}$$

- (b) if the particles are fermions? If the particles are fermions, there is only one state $\psi_a(1)\psi_b(2)\psi_c(3)$.]

Formulae

$$-\frac{\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial x^2} + V\Psi = i\hbar \frac{\partial \Psi}{\partial t}$$

$$H\Psi = i\hbar \frac{\partial \Psi}{\partial t}$$

$$\hat{p} = -i\hbar \frac{\partial}{\partial x}$$

$$[x, p] = i\hbar$$

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}; \quad \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}; \quad \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

- Time independent perturbation theory

$$E_n^1 = \langle \psi_n^0 | H' | \psi_n^0 \rangle$$

$$\psi_n^1 = \sum_{m \neq n} \frac{\langle \psi_m^0 | H' | \psi_n^0 \rangle}{(E_n^0 - E_m^0)} \psi_m^0$$

$$E_n^2 = \sum_{m \neq n} \frac{|\langle \psi_m^0 | H' | \psi_n^0 \rangle|^2}{(E_n^0 - E_m^0)}$$

- One dimensional infinite square well

$$\psi_n(x) = \left(\frac{2}{a}\right)^{1/2} \sin\left(\frac{n\pi}{a}x\right)$$

$$E_n = \frac{\pi^2 \hbar^2}{2ma^2} (n^2)$$

- Virial Theorem

$$2\langle T \rangle = \langle \mathbf{r} \cdot \nabla V \rangle$$

- Generators

$$e^{i(\boldsymbol{\sigma} \cdot \hat{n})\phi/2} = \cos(\phi/2) + i(\hat{n} \cdot \boldsymbol{\sigma}) \sin(\phi/2)$$

- Integral

$$\int_0^a \sin^2 \frac{n\pi x}{a} dx = \frac{a}{2}$$