

EXPANSION OF MAGNETIC FIELD IN CYLINDRICAL COORDINATES.

D. RUBIN

In order to incorporate measurements of the azimuthal dependence of vertical and radial fields into a tracking program, we

- Write the general solution to Laplace's equation in cylindrical coordinates, $\nabla^2\psi(\rho, \theta, z) = 0$
- Compute $\nabla\psi = \mathbf{B}$
- Fit $B_\rho(\rho = \rho_0, \theta, z = 0)$ to the measured radial field data to determine coefficients
- Fit $B_z(\rho = \rho_0, \theta, z = 0)$ to the measured vertical field data the rest of the coefficients

The general solution is

$$\psi = \sum_{m=0}^{\infty} \sum_{n=1}^{\infty} J_m(k_{mn}\rho) (\sinh(k_{mn}z)(A_{mn} \sin m\phi + B_{mn} \cos m\phi) + \cosh(k_{mn}z)(C_{mn} \sin m\phi + D_{mn} \cos m\phi))$$

Then

$$\begin{aligned} B_z &= -\frac{\partial\psi}{\partial z} \\ &= \sum_{m=0}^{\infty} \sum_{n=1}^{\infty} J_m(k_{mn}\rho) k_{mn} (\cosh(k_{mn}z)(A_{mn} \sin m\phi + B_{mn} \cos m\phi) \\ &\quad - \sinh(k_{mn}z)(C_{mn} \sin m\phi + D_{mn} \cos m\phi)) \end{aligned}$$

At $\rho = \rho_0, z = 0,$

$$B_z(\rho_0, \phi, 0) = \sum_{m=0}^{\infty} \sum_{n=1}^{\infty} J_m(k_{mn}\rho_0) k_{mn} (A_{mn} \sin m\phi + B_{mn} \cos m\phi)$$

$$\begin{aligned} B_\rho &= -\frac{\partial\psi}{\partial\rho} \\ &= \sum_{m=0}^{\infty} \sum_{n=1}^{\infty} \frac{\partial J_m}{\partial\rho}(k_{mn}\rho) (\sinh(k_{mn}z)(A_{mn} \sin m\phi + B_{mn} \cos m\phi) \\ &\quad + \cosh(k_{mn}z)(C_{mn} \sin m\phi + D_{mn} \cos m\phi)) \end{aligned}$$

$$B_\rho(\rho_0, \phi, 0) = \sum_{m=0}^{\infty} \sum_{n=1}^{\infty} \frac{\partial J_m}{\partial\rho}(k_{mn}\rho_0) (C_{mn} \sin m\phi + D_{mn} \cos m\phi)$$

Write the measured fields along the magic radius

$$B_z^m = \sum_m (a_m \cos m\phi + b_m \cos m\phi)$$

$$B_\rho^m = \sum_m (c_m \cos m\phi + d_m \cos m\phi)$$

Then

$$a_m = \sum_{n=1}^{\infty} J_m(k_{mn}\rho_0) k_{mn} A_{mn}$$

$$b_m = \sum_{n=1}^{\infty} J_m(k_{mn}\rho_0) k_{mn} B_{mn}$$

$$c_m = \sum_{n=1}^{\infty} \frac{\partial J_m}{\partial \rho}(k_{mn}\rho_0) k_{mn} C_{mn}$$

$$d_m = \sum_{n=1}^{\infty} \frac{\partial J_m}{\partial \rho}(k_{mn}\rho_0) k_{mn} D_{mn}$$

Define

$$A'_{mn} = J_m(k_{mn}\rho_0) k_{mn} A_{mn}$$

$$B'_{mn} = J_m(k_{mn}\rho_0) k_{mn} B_{mn}$$

$$C'_{mn} = \frac{\partial J_m}{\partial \rho}(k_{mn}\rho_0) k_{mn} C_{mn}$$

$$D'_{mn} = \frac{\partial J_m}{\partial \rho}(k_{mn}\rho_0) k_{mn} D_{mn}$$

And

$$\sum_n A'_{mn} = a_m$$

etc. It appears that since I have no information at nonzero z or $\rho \neq \rho_0$ that all of the k_{mn} are degenerate. So I might as well only include $n = 1$ in the sum. The result is a perfectly good solution although perhaps not the most general. At least it will obey Maxwell and

be consistent with the boundary conditions. Then

$$\begin{aligned} a_m &= J_m(k_{m1}\rho_0)k_{m1}A_{m1} \\ b_m &= J_m(k_{m1}\rho_0)k_{m1}B_{m1} \\ c_m &= \frac{\partial J_m}{\partial \rho}(k_{m1}\rho_0)k_{m1}C_{m1} \\ d_m &= \frac{\partial J_m}{\partial \rho}(k_{m1}\rho_0)k_{m1}D_{m1} \end{aligned}$$

Let's drop the subscript n

Assume that $|\mathbf{B}| = \mathbf{B}_z$. a_m, b_m and c_m, d_m are the coefficients of the fits to the measured B_z and B_r respectively. Define

$$\begin{aligned} A_m &= a_m/(J_m(k_{m1}\rho_0)k_{m1}) \\ B_m &= b_m/(J_m(k_{m1}\rho_0)k_{m1}) \\ C_m &= c_m/(\frac{\partial J_m}{\partial \rho}(k_{m1}\rho_0)k_{m1}) \\ D_m &= d_m/(\frac{\partial J_m}{\partial \rho}(k_{m1}\rho_0)k_{m1}) \end{aligned}$$

Then the fields

$$\begin{aligned} B_z &= \sum_{m=0}^{\infty} J_m(k_m\rho)k_m (\cosh(k_m z)(A_m \sin m\phi + B_m \cos m\phi) \\ &\quad - \sinh(k_m z)(C_m \sin m\phi + D_m \cos m\phi)) \\ B_\rho &= \sum_{m=0}^{\infty} \frac{\partial J_m}{\partial \rho}(k_m\rho) (\sinh(k_m z)(A_m \sin m\phi + B_m \cos m\phi) \\ &\quad + \cosh(k_m z)(C_m \sin m\phi + D_m \cos m\phi)) \\ B_\phi &= \sum_{m=0}^{\infty} \frac{J_m(k_m\rho)}{\rho} m (\sinh(k_m z)(A_m \cos m\phi - B_m \sin m\phi) \\ &\quad + \cosh(k_m z)(C_m \cos m\phi - D_m \sin m\phi)) \end{aligned}$$