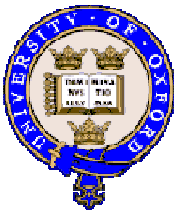

The Charm Renaissance: D-physics – a Selective Review

- Charm physics: why we live in interesting times
- Experimental facilities and techniques
- Charm and precision CKM physics
- Searches for new physics in charm decays
- Conclusions



Guy Wilkinson (University of Oxford)
On behalf of the CLEO-c collaboration



Charm physics: extinct, dormant or active ?



A common view:

“Charm physics: illustrious past which played essential role in foundation of Standard Model, but not much present & still less future. B physics is where it’s all at.”

Some reasons that charm may appear underwhelming:

- Slow oscillation rate;
- Very low level of CP violation;
- Interesting physics masked by long distance dynamics.

Is this fair, or does charm remain a very important research area, indeed one where new interest is erupting right now?



A Sporting Analogy

The future is bright (cf. b physics)...



English and European champions

Yesterday's men (cf. c physics):



18 times league championships –
but none for 18 years (and counting)

Charm physics – a new dawn

There are several good reasons why charm physics is back in the limelight:

1) Precision CKM tests

Success of the B-factories and the Tevatron has meant that unitarity triangle tests are entering a new, precision era. Although the CKM elements being studied are those accessible in B-decays, charm turns out to be a vital ingredient in programme.

2) Charm mixing and its legacy

Discovery of D^0 - \bar{D}^0 oscillation was HEP highlight of last 2 years. Higher than expected rate is (arguably) intriguing in its own right, and points the way forward to searches for CPV.

3) Recent discoveries in spectroscopy

Discovery of narrow c - \bar{c} states above threshold (X,Y,Z) and narrow excited D_s states. (But not discussed further today.)

Charm physics certain to have an exciting few years ahead (unlike Liverpool FC)

Facilities, experimental attributes and Dalitz plots

Charm Facilities – Recent Past, Present & Future

Fixed target experiments
eg. E687, E791, FOCUS

CLEO,
B-factories

Very high statistics from e^+e^- continuum in all final states. Highlight (so far): mixing discovery

CLEO-c

Threshold running: double tags and quantum coherence

BES-III

(~20x CLEO dataset?)

Tevatron, esp. CDF

First hadron collider D decay studies with prompt charm (eg. mixing: PRL 100 (2008) 121802)

LHCb

Charm from B's in selected final states (~ 10^2 higher statistics than B-factory)

Super-flavour factory

10-100x B-factory statistics. D's from e^+e^- continuum *and* in dedicated threshold runs

Super-LHCb

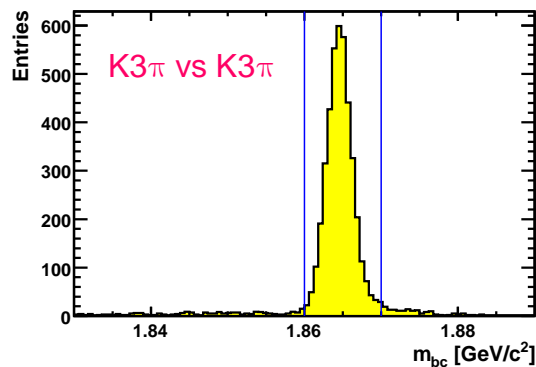
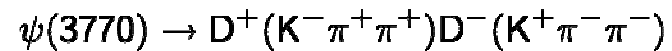
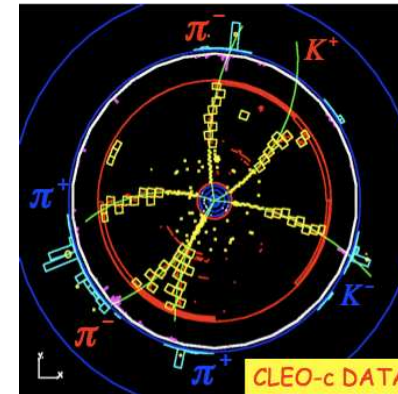
10x luminosity and more efficient + inclusive trigger

Charm Studies at Threshold (CLEO-c)

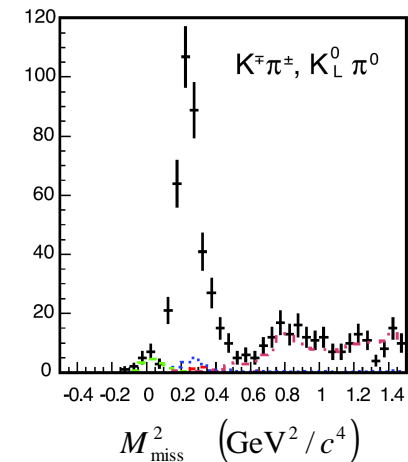
CLEO-c accumulated 818 pb⁻¹ at $\psi(3770)$
and 586 pb⁻¹ at 4170 MeV (for D_s production)

Some advantages of threshold running:

- Very clean – no fragmentation particles.
Double tag studies have v. low bckgds



- Unseen particle reconstruction through kinematic constraints, eg. K_L⁰ and ν .

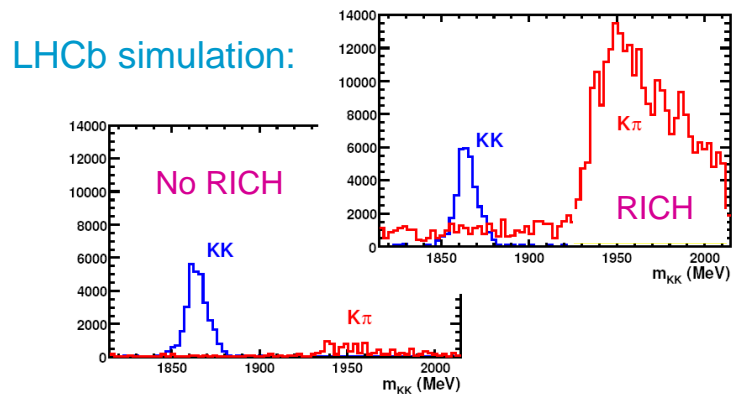


- Quantum coherence. $\psi(3770)$ gives C=-1 state → CP-tagging possible !

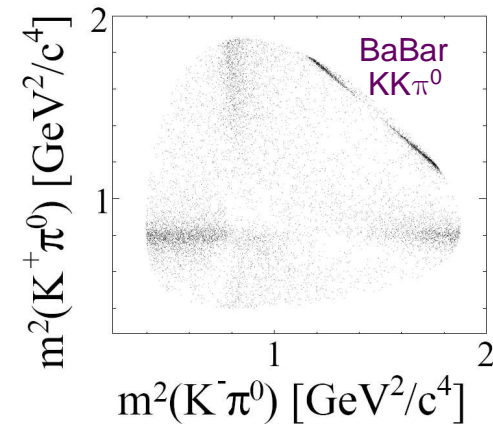
Necessary Experimental Attributes

Characteristics needed for successful D physics ~ those required for B physics

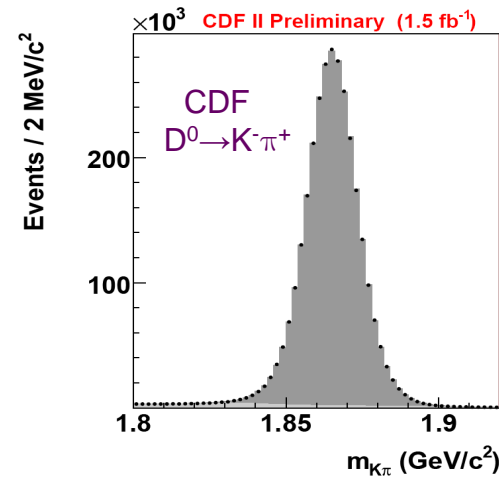
- Efficient tracking and (if possible) good calorimetry for γ and π^0 reconstruction
- Hadron identification abilities



- In hadronic environment need trigger system sensitive to final states of interest



PRD 76 (2007) 011102 (R)



PRL 100 (2008) 121802

Experimental Techniques: Dalitz Plots

Dalitz plot is invaluable technique exploited in many charm analyses.



Richard Dalitz
1925-2006

Kinematics of 3-body decay $D \rightarrow A, B, C$ fully described by 2 parameters. Typical choice:

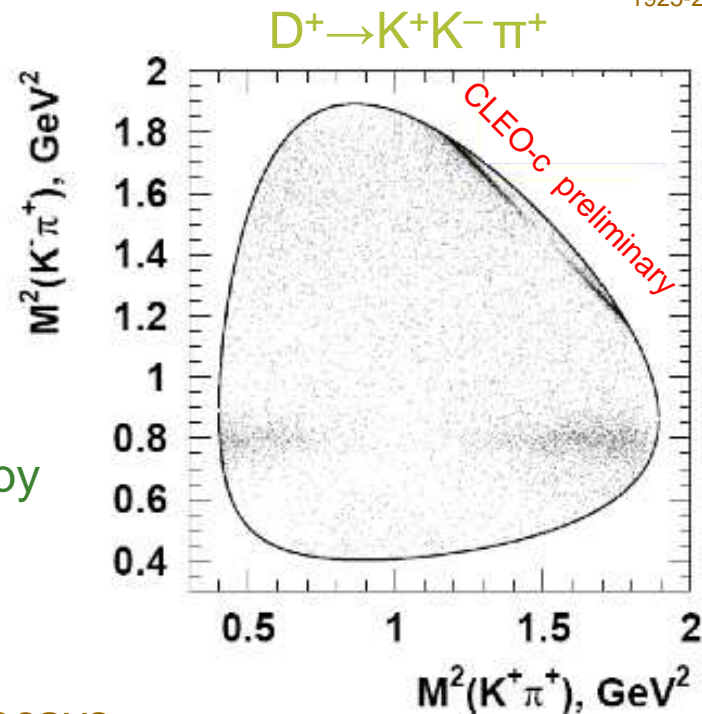
$$m_{AB}^2 \equiv (p_A + p_B)^2 ; m_{BC}^2 \equiv (p_B + p_C)^2$$

Lorentz invariant, and phase space flat.
Allows resonances to be clearly seen.

Charm Dalitz plots have many uses:

- 1) as a probe of light meson spectroscopy
- 2) key role in the CKM- γ measurement
- 3) mixing and CPV studies

Can be extended to 4- (and more) body decays



Dalitz Plots and Resonance Models

To extract physics contributing to a Dalitz plot, necessary to develop a model
 Common choice: isobar model - fit a_j, b, α_j, β :

$$A(m_{12}^2, m_{13}^2) = \left[\sum_j a_j e^{i\alpha_j} A_j(m_{12}^2, m_{13}^2) \right] + b e^{i\beta}$$

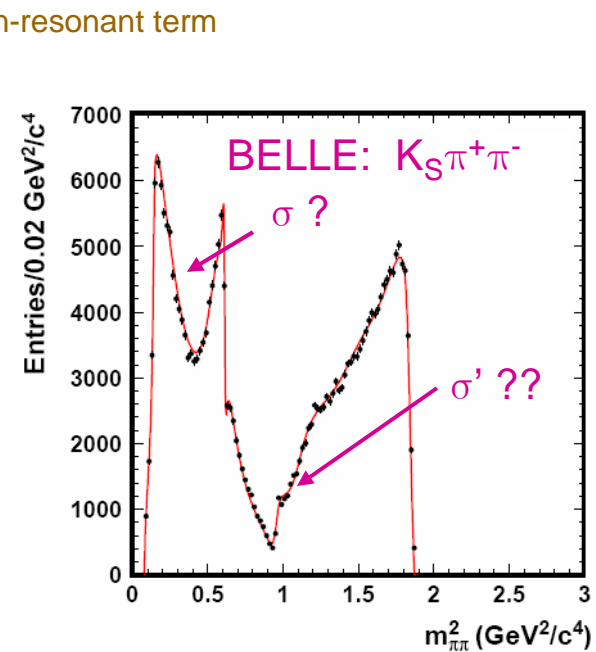
↑ Sum over resonances
 ↑ Non-resonant term

where amplitudes described as follows:

$$A_j(m_{12}^2, m_{13}^2) = F_D^J F_D^r \cdot M_r^j \cdot BW_r^j$$

↑ Blatt-Weisskopf form factors
 (angular momentum barrier penetration)
 ↑ Relativistic Breit-Wigner

↑ Angular distribution



Works well for P-, D-wave, but not so well for S-wave.

Not great for broad overlapping resonances, eg. σ, σ', κ ? (& unitarity violated).

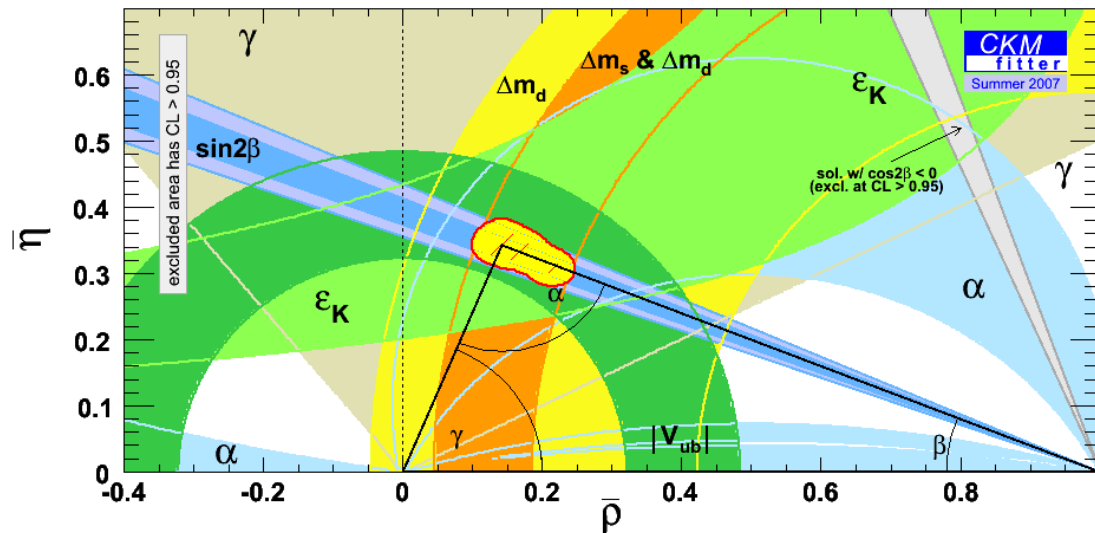
Here alternative treatments preferred based on scattering data, eg. LASS for $K\pi$ (Aston et al. Nucl. Phys. B 296 (1988) 493), and K-matrix for $\pi\pi$ (I.J.R.Aitchison, Nucl. Phys. A189 (1972) 417)

D decays and the CKM unitarity triangle

- Measurements of γ with $B \rightarrow DK$
- Lattice QCD tests and the ‘mixing side’
- (Not covered: semileptonic decay rates and form factors, branching ratio measurements etc)

The Unitarity Triangle and D decays

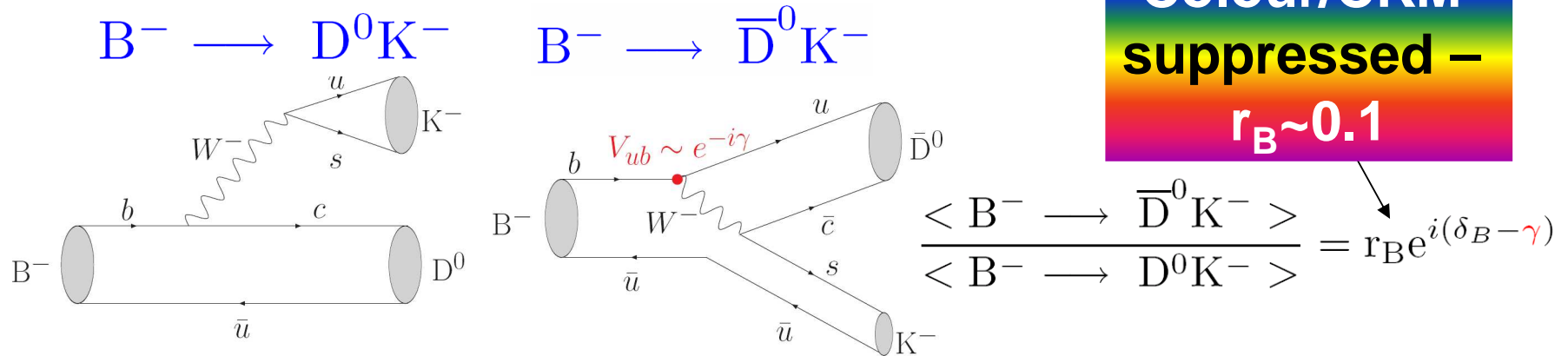
Classical unitarity triangle is constrained by quantities measured in B decays



But key measurements have high dependence (direct & indirect) on D decays:

- Tree-level measurements of angle γ
At present $\gamma^{\text{measured}} = 77 \pm 31^\circ$
 - Lattice QCD input to side opposite to γ
This side and $\sin 2\beta$ gives $\gamma^{\text{predicted}} = 68 \pm 4^\circ$
- ← Improve this, verify that, and compare...

γ from $B^\pm \rightarrow DK^\pm$

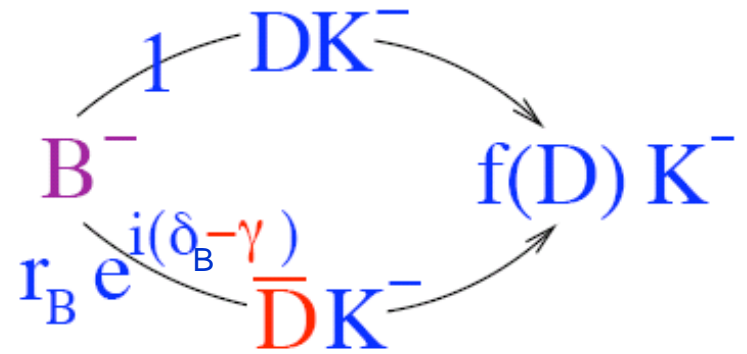


- Extraction through interference between $b \rightarrow u$ and $b \rightarrow c$ transitions

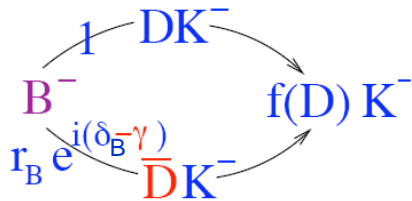
- Require D^0 and \bar{D}^0 decay to a common final state, $f(D)$. Today consider 2 possibilities:

$K_S^0 \pi \pi$; $K \pi$; ($K \pi \pi \pi$ - see backups)

- Tree level processes: little sensitivity to New Physics \rightarrow SM ‘standard candle’



Dalitz Plots for γ at Belle & BaBar

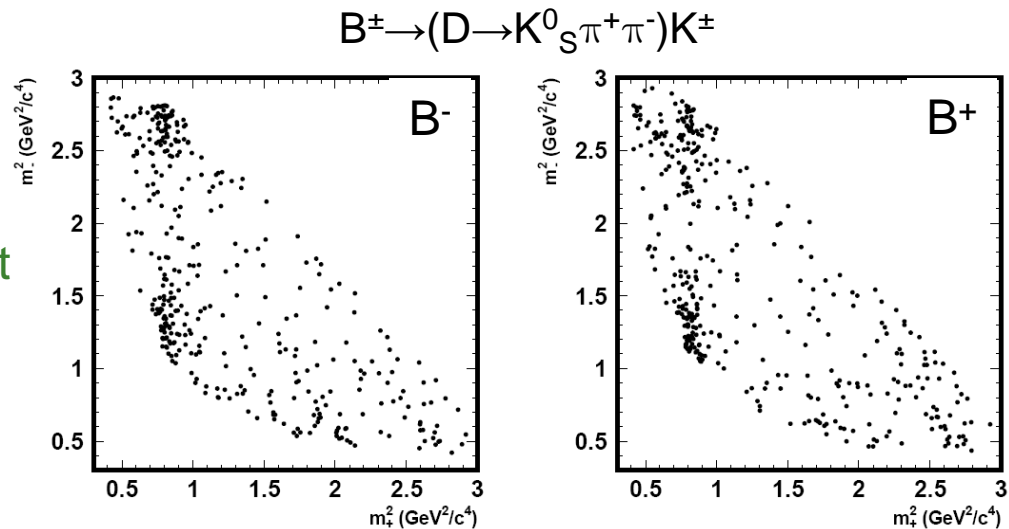


A powerful (and at *present*, only statistically useful) choice of common state $f(D)$ is $K_S \pi^+ \pi^-$. Rich resonant substructure.

Differences between B^- and B^+ Dalitz plots allow γ to be extracted in unbinned fit...

...need to understand different amplitudes from D^0 and \bar{D}^0 decay across Dalitz space, esp. variation in strong phase

Need a D decay model !



BELLE: arXiv:0803.3375

BaBar (383M BB) $\gamma = 76^\circ \pm 22^\circ(\text{stat}) \pm 5^\circ(\text{sys}) \pm 5^\circ(\text{model})$ $r_B = 0.09 \pm 0.09$

BELLE (657M BB) $\gamma = 76^\circ_{-13^\circ}^{+12^\circ}(\text{stat}) \pm 4^\circ(\text{sys}) \pm 9^\circ(\text{model})$ $r_B = 0.16 \pm 0.04$

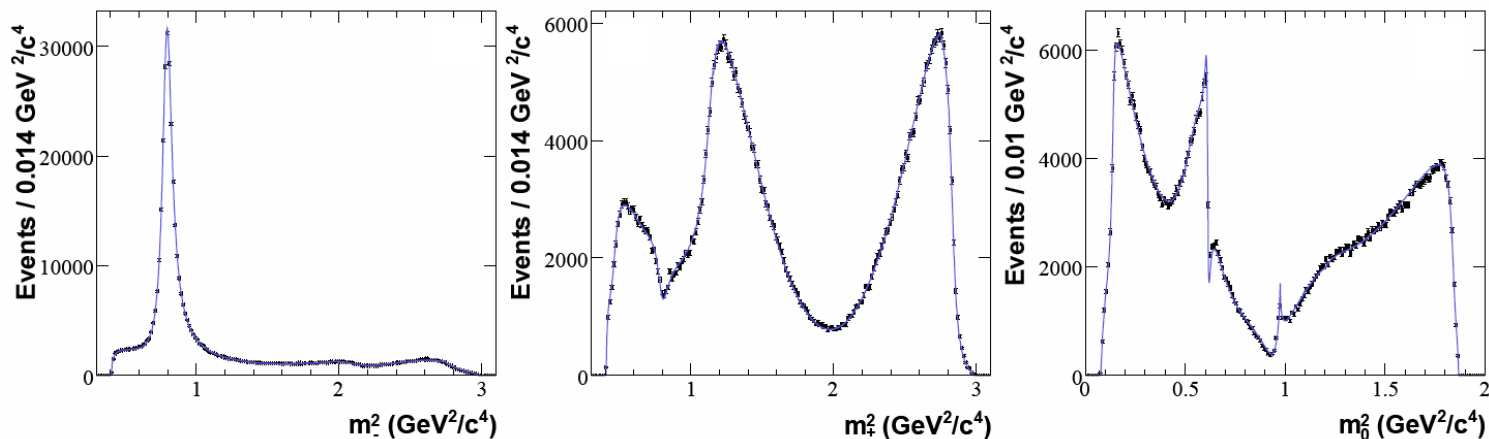
LHCb with 10 fb^{-1} can approach 3° statistical error

(NB sensitivity scales with r_B . All $B \rightarrow DK$ data suggest ~ 0.10)

Modelling the $K_s \pi^+ \pi^-$ decay

Unbinned fit of Dalitz space in $B \rightarrow D(K_s \pi^+ \pi^-) K$ decays requires reliable model of D decay. Model developed on flavour tagged D^* decays.

State of the art – BaBar model fitted from 487k decays:



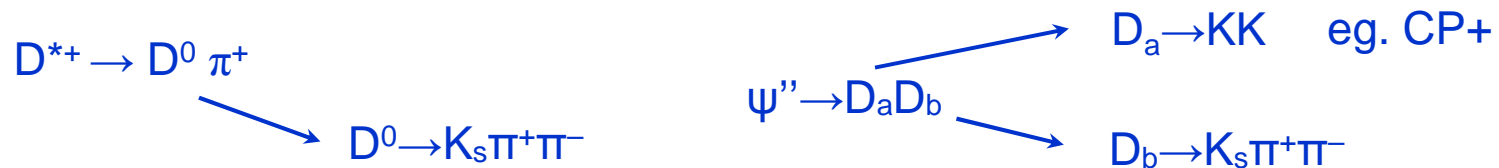
arXiv:0804.2089

Ingredients – 10 resonances described with isobar model. S-wave $\pi\pi$ and $K\pi$ treated with K-matrix approach and LASS parametrisation respectively ($\chi^2 / \text{ndf} = 1.11$ to be compared with 1.20 for pure isobar model)

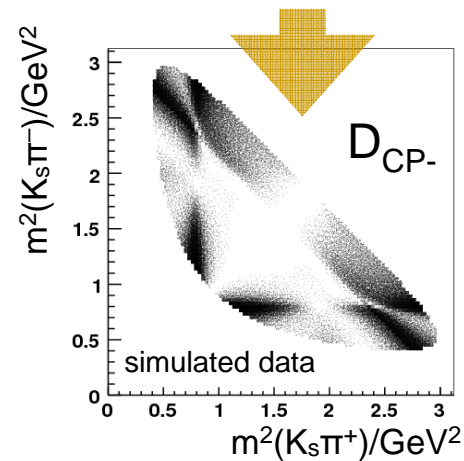
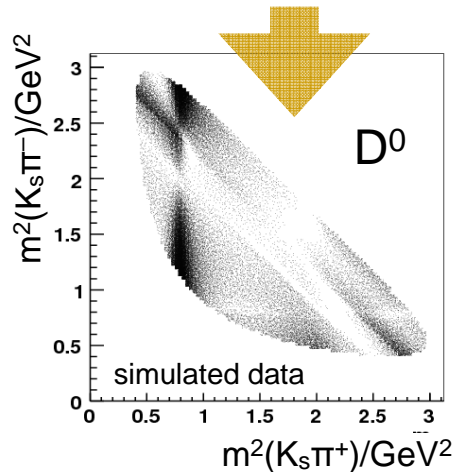
Impressive work – error on γ estimated to be 5° . But model systematic, even this small, uncomfortable for future very high stats measurements eg. LHCb.

CP- tagged Dalitz Plots

Dalitz plots of CP-tagged decays at the Ψ'' provide orthogonal info to flavour tagged events accessible in, eg., D^* decays. They access the strong phase difference between the D^0 & \bar{D}^0 – vital information in the γ measurement



Flavour tagged distribution $\propto |D^0|^2$ or $|\bar{D}^0|^2$



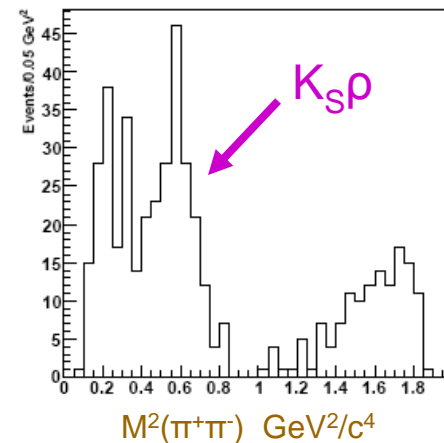
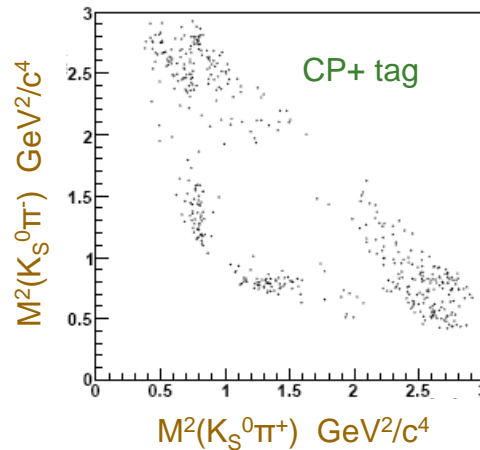
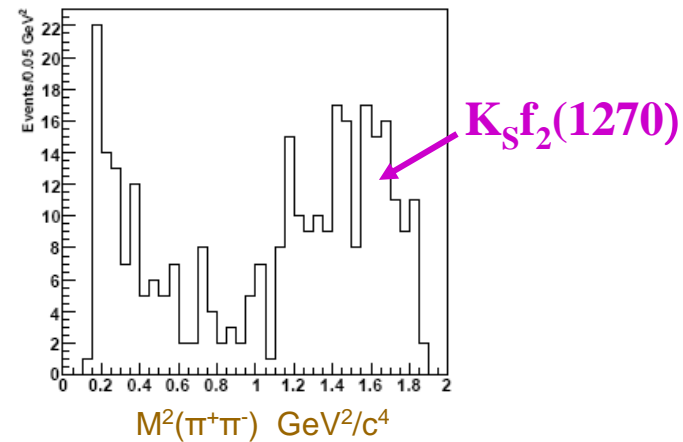
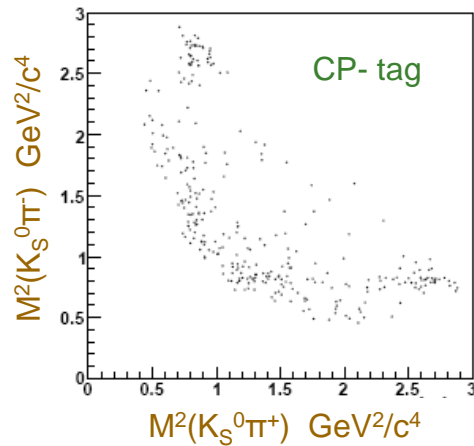
CP-tagged $\propto |D^0|^2 + |\bar{D}^0|^2 \pm 2 |D^0 \bar{D}^0| \cos \delta$

With both flavour and CP-tagged data we may either validate model, or avoid model *entirely* and used measured quantities as input to binned fit !

CP-tagged $K_S \pi^+ \pi^-$ Dalitz plots

Clear differences seen between CP-odd and CP-even:

CLEO-c preliminary,
~800 CP-tagged $K_S \pi^+ \pi^-$



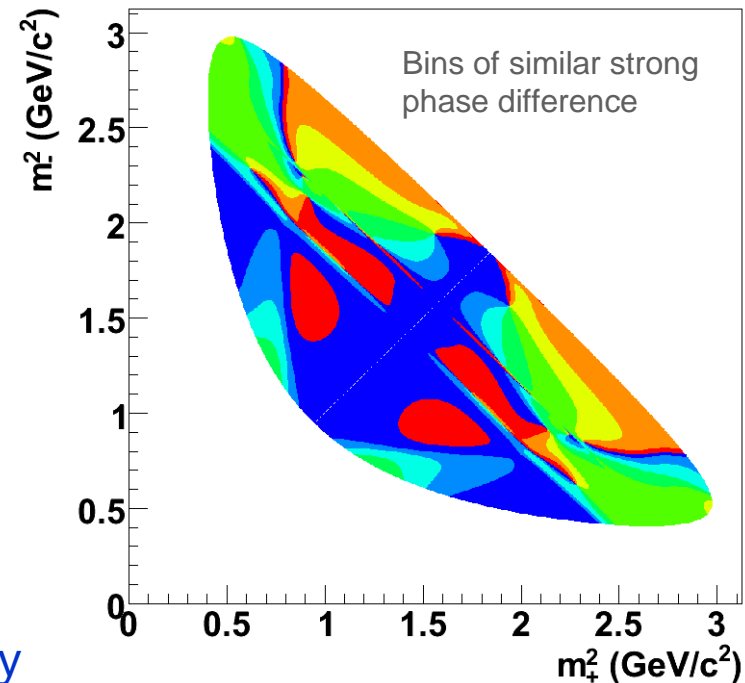
Binned Analysis of $B^\pm \rightarrow D(K_s \pi^+ \pi^-) K^\pm$

CP-tagged data can be used to avoid entirely need for model.

Expected number of events in bins of Dalitz space can be expressed in terms of γ , r_B , δ_B and measured yields in CP-tagged D decays!

(Additional input comes from quantum correlated $K_s \pi \pi$ vs $K_s \pi \pi$ events)

Choice of bins informed by model, in order to maximise statistical sensitivity (only ~20% degradation w.r.t. unbinned fit)



Bondar and Poluektov, arXiv:0801.0840

CLEO-c analysis in progress...

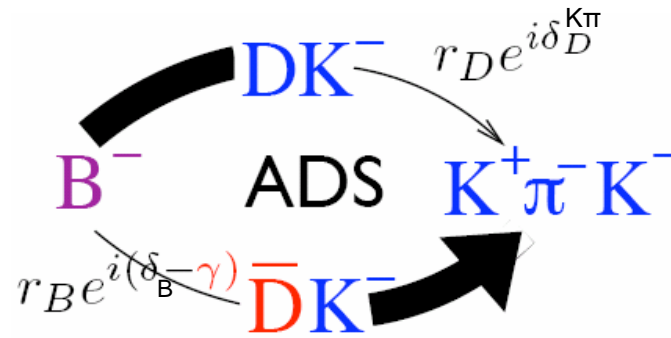
But no model error! Only residual uncertainty from finite CLEO-c statistics ($\sim 3^\circ$)

Atwood-Dunietz-Soni (ADS) Method

Low interference scale of $B \rightarrow DK$ method ($r_B \sim 0.1$) can be enhanced by exploiting Doubly Cabibbo Suppressed modes eg. $D^0 \rightarrow K^+\pi^-$

This introduces two new parameters:

$$\frac{\langle D^0 \rightarrow K^+\pi^- \rangle}{\langle \bar{D}^0 \rightarrow K^+\pi^- \rangle} = r_D^{K\pi} e^{i\delta_D^{K\pi}}$$



$r_D^{K\pi}$ known well, $\delta_D^{K\pi}$ unknown

~ 0.06 , ie. similar in magnitude to r_B

4 possible final states, between 2 of which there can be a big CP-asymmetry:

$$\Gamma(B^- \rightarrow (K^+\pi^-)_D K^-) \propto r_B^2 + (r_D^{K\pi})^2 + 2r_B r_D^{K\pi} \cdot \cos(\delta_B + \delta_D^{K\pi} - \gamma)$$

$$\Gamma(B^+ \rightarrow (K^-\pi^+)_D K^+) \propto r_B^2 + (r_D^{K\pi})^2 + 2r_B r_D^{K\pi} \cdot \cos(\delta_B + \delta_D^{K\pi} + \gamma)$$

A powerful way to constrain γ , but need to know $\delta_D^{K\pi}$
 Can be measured in quantum correlated D decays !

these interference terms are 1st order

Measuring $\delta_D^{K\pi}$ in Quantum Correlated D Decays

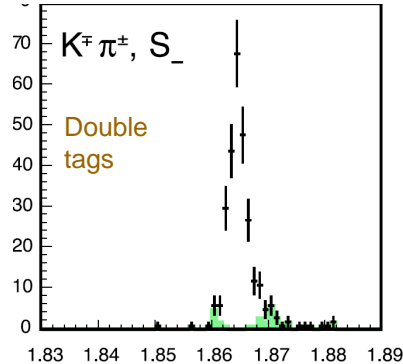
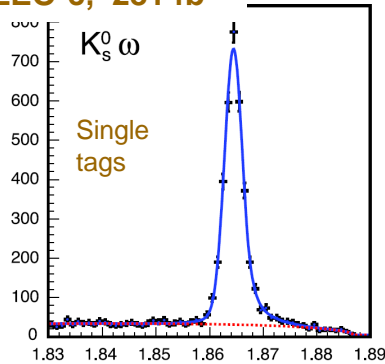
Idea: tag one D in CP eigenstate, other side is mixture of D^0 and \bar{D}^0 , hence:

$$\text{Rate} \sim B_{CP+} B_{K\pi} (1 + 2r_D^{K\pi} \cos \delta_D^{K\pi})$$

Approximate - full expression has additional dependence on mixing parameters x & y ...

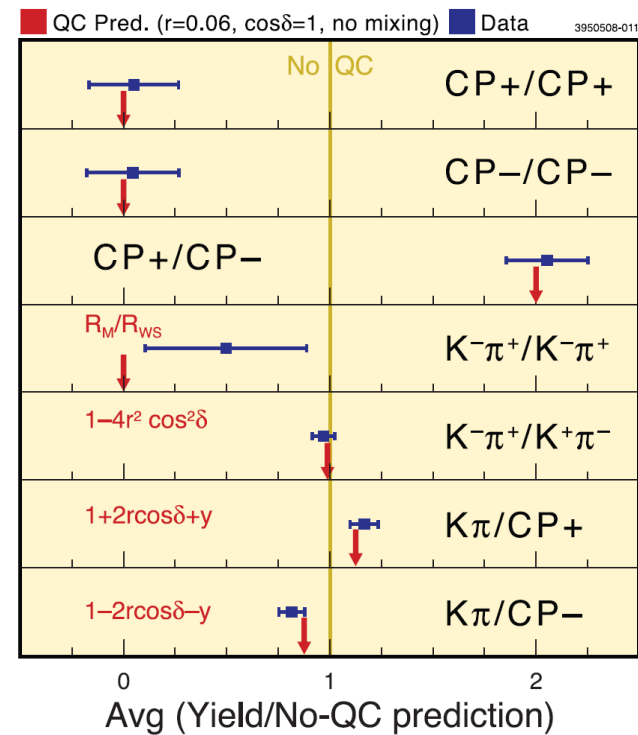
Analysis: measure set of single & double tag rates, with $K\pi$ vs CP tags, & flavour tags

CLEO-c, 281 fb⁻¹



Mass / GeV²/c⁴

Extract $\delta_D^{K\pi}$, plus results on other parameters, including branching ratios.



Quantum correlations clearly seen !

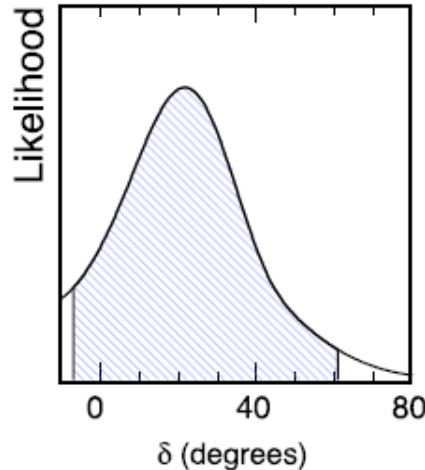
CLEO-c: PRL 100 (2008) 221801; arXiv:0802.2268

CLEO-c 281 fb⁻¹ Results for $\delta_D^{K\pi}$

Result also important for charm mixing
(x' , y' measured in 'wrong sign' $K\pi$
analysis related to x , y through:
 $x' = x \cos \delta_D^{K\pi} + y \sin \delta_D^{K\pi}$)

Most precise result
on $\delta_D^{K\pi}$ obtained
with mixing results
used as external
constraint:

$$\delta_D^{K\pi} = \left(22_{-12-11}^{+11+9} \right)^\circ$$



Fit results with all external constraints

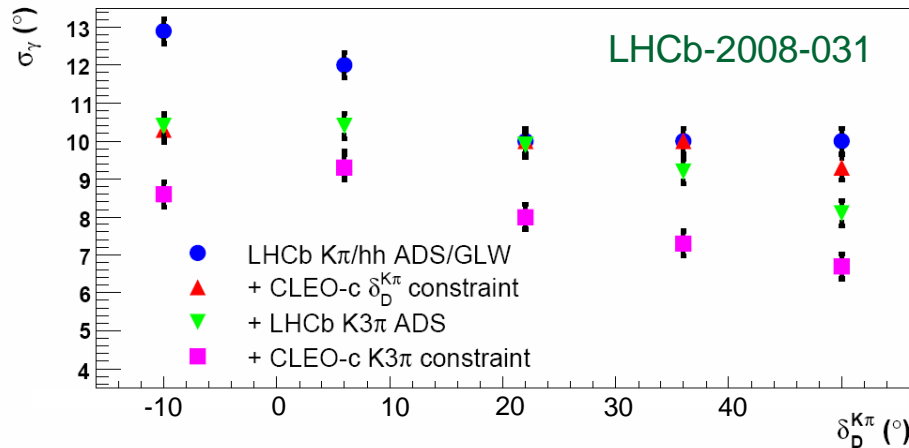
Parameter	Extended Fit
\mathcal{N} (10^6)	$1.042 \pm 0.021 \pm 0.010$
y (10^{-3})	$6.5 \pm 0.2 \pm 2.1$
r^2 (10^{-3})	$3.44 \pm 0.01 \pm 0.09$
$\cos \delta$	$1.10 \pm 0.35 \pm 0.07$
x^2 (10^{-3})	$0.06 \pm 0.01 \pm 0.05$
$x \sin \delta$ (10^{-3})	$4.4 \pm 2.4 \pm 2.9$
$K^- \pi^+$ (%)	$3.78 \pm 0.05 \pm 0.05$
$K^- K^+$ (10^{-3})	$3.88 \pm 0.06 \pm 0.06$
$\pi^- \pi^+$ (10^{-3})	$1.36 \pm 0.02 \pm 0.03$
$K_S^0 \pi^0 \pi^0$ (10^{-3})	$8.35 \pm 0.44 \pm 0.42$
$K_S^0 \pi^0$ (%)	$1.14 \pm 0.03 \pm 0.03$
$K_S^0 \eta$ (10^{-3})	$4.42 \pm 0.15 \pm 0.28$
$K_S^0 \omega$ (%)	$1.12 \pm 0.04 \pm 0.05$
$X^- e^+ \nu_e$ (%)	$6.59 \pm 0.16 \pm 0.16$
$K_L^0 \pi^0$ (%)	$1.01 \pm 0.03 \pm 0.02$
$\chi_{\text{fit}}^2/\text{ndof}$	$55.3/57$

Result will improve with full 818 fb⁻¹ data set,
use of additional tags, and possible exploitation of 4170 MeV data.

Similar analysis can be done for $K\pi\pi\pi$ – but here need to worry about resonant substructure which brings dilution to ADS interference (see backup slides).

Expected Sensitivity to γ at LHCb

Expected γ precision with 2 fb^{-1} of data (one year) for ADS modes alone:



Improvements in going from 2-body ADS to 4-body ADS & adding constraints from CLEO-c

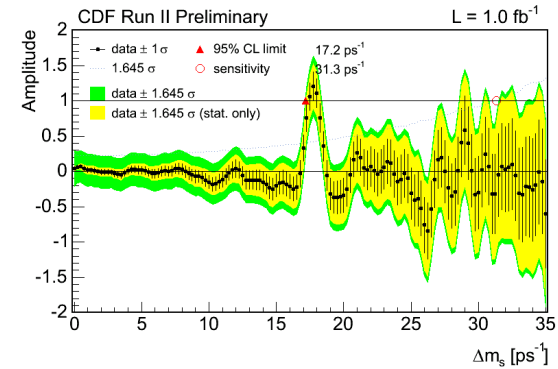
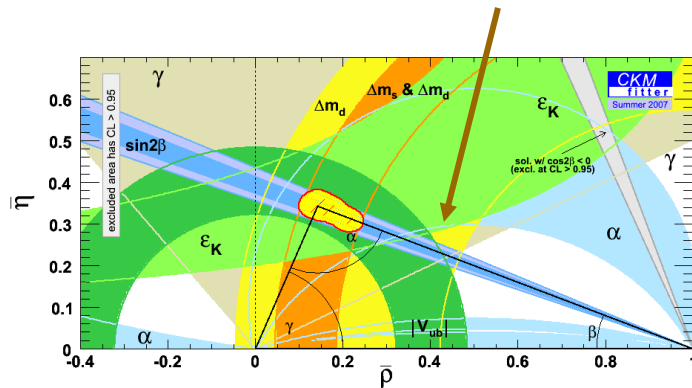
Add other measurements, especially $B \rightarrow D(K_S \pi^+ \pi^-) K$, & extrapolate to 10 fb^{-1}

$\sigma_\gamma = 1.9 - 2.7^\circ$...in which $B \rightarrow DK$ methods have a weight of $\sim 70\%$ (variation in number depends on values of phases)

Understanding of D decay properties central to precise measurement of γ !

Lattice QCD and the ‘Mixing Side’

‘Mixing side’ of unitarity triangle determined by B_d/B_s oscillation rates – box diagrams sensitive to new physics – and QCD corrections



Very well known (~0.3%), since observation of B_s mixing

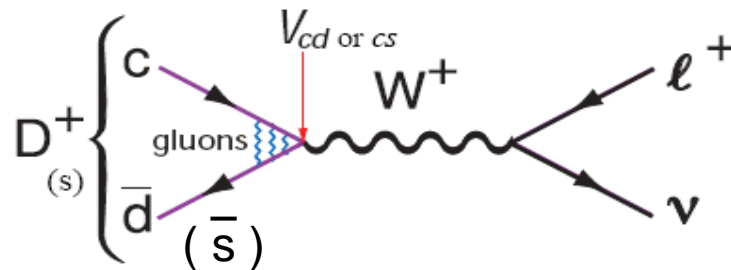
$$\text{Length} = (f_{B_s} \sqrt{B_{B_s}}) / (f_{B_d} \sqrt{B_{B_d}}) \sqrt{(\Delta m_d m_{B_d}) / (\Delta m_s m_{B_s})} = |V_{td}/V_{ts}|$$

Calculated on lattice – present *assigned* uncertainty ~5% (and will decrease)

Highly desirable to cross-check lattice against experiment - go to D system !

Leptonic D Decays and Decay Constants

In D^+ and D_s c and spectator quark can annihilate to produce leptonic final state:



In general, for all pseudoscalars:

$$\Gamma(P^+ \rightarrow \ell^+ \nu) = \frac{1}{8\pi} G_F^2 f_P^2 m_\ell^2 M_P \left(1 - \frac{m_\ell^2}{M_P^2}\right)^2 |V_{q_1}|^2$$

Since V_{cd} and V_{cs} well known, can extract f_D and f_{D_s} and compare with lattice !

Measurements of $D_{(s)} \rightarrow l\nu$ Branching Fractions

Precise measurements now exist for:

$\mu^+\nu$, $\tau^+ (\rightarrow \pi^+\nu)\nu$ CLEO-c (PRL 99 (2007) 071802; arXiv:0704.0437 + prelim)

D_s $\mu^+\nu$ BELLE (arXiv:0709.1340) & BaBar (hep-ex/0607094)

$\tau^+ \rightarrow (e^+\nu\nu)\nu$ CLEO-c (PRL 100 (2008) 161801)

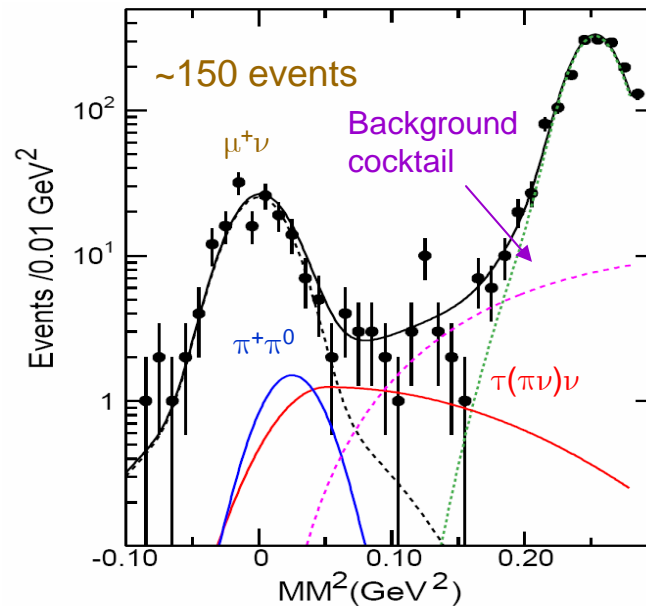
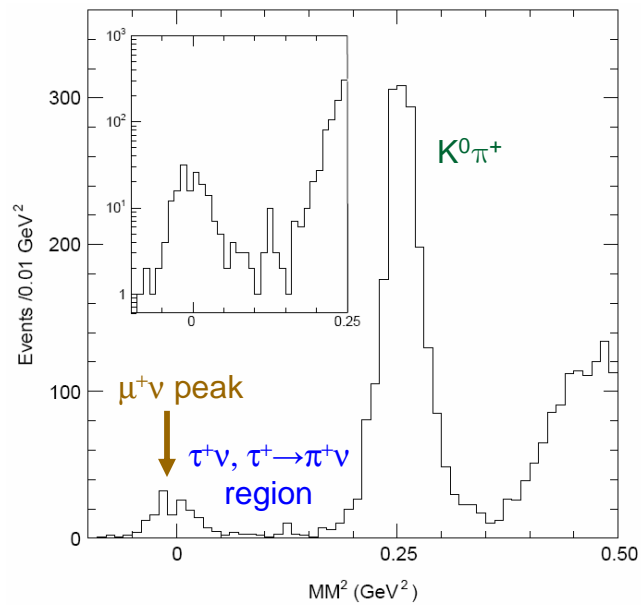
D^+ CLEO-c (arXiv:0806.2112)

Basic method for $\mu\nu$ measurement:

- CLEO-c: for f_D reconstruct one D^+ , look for MIP (μ), and then compute missing mass squared (similar for f_{D_s} , but here exploit $D_s D_s^*$ production in 4170 MeV dataset)
- B-factory: infer presence of D_s from recoiling mass against reconstructed D & fragmentation. Add candidate μ and compute missing mass

CLEO-c $D^+ \rightarrow \mu^+ \nu$

Missing mass squared distribution (incl. log zoom with fit):



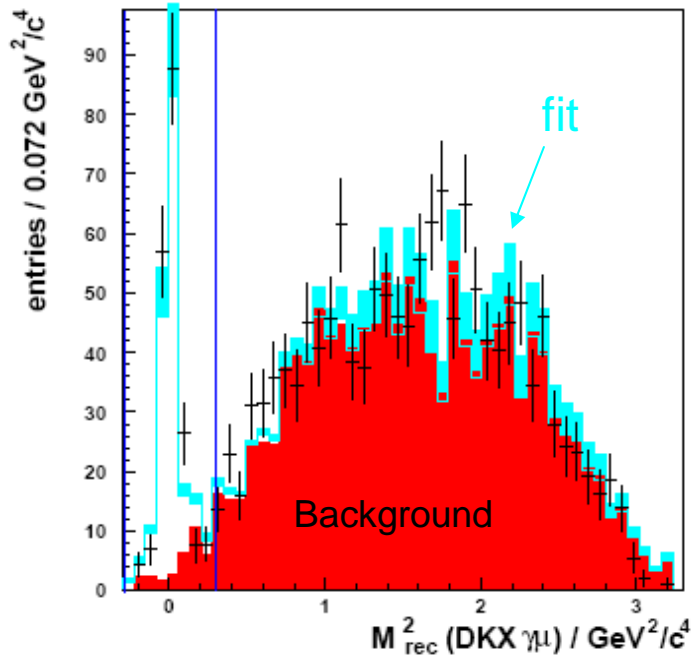
$$\text{BR}(D^+ \rightarrow \mu^+ \nu) = (3.82 \pm 0.32 \pm 0.09) \times 10^{-4}$$

$$f_D = (205.8 \pm 8.5 \pm 2.5) \text{ MeV}$$

(result with τ_{ν}/μ_{ν} fixed
at SM expectation)

$D_s \rightarrow \mu^+ \nu$

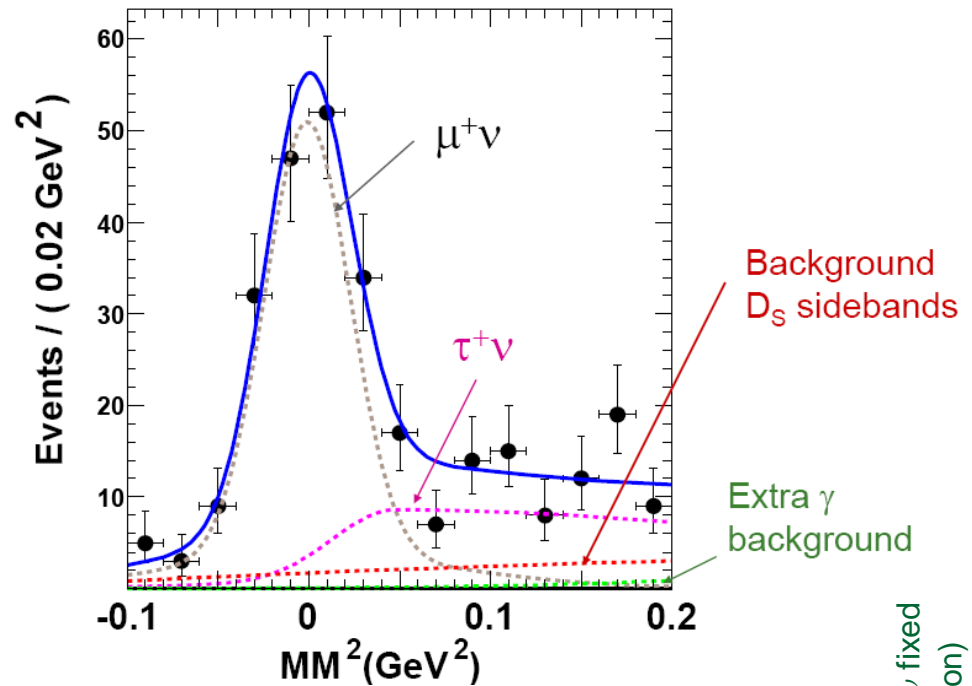
Belle: 548 fb⁻¹



$$\text{BR}(D_s \rightarrow \mu^+ \nu) = (6.44 \pm 0.76 \pm 0.57) \times 10^{-3}$$

$$f_{D_s} = (275 \pm 16 \pm 12) \text{ MeV}$$

CLEO-c preliminary: 424 pb⁻¹



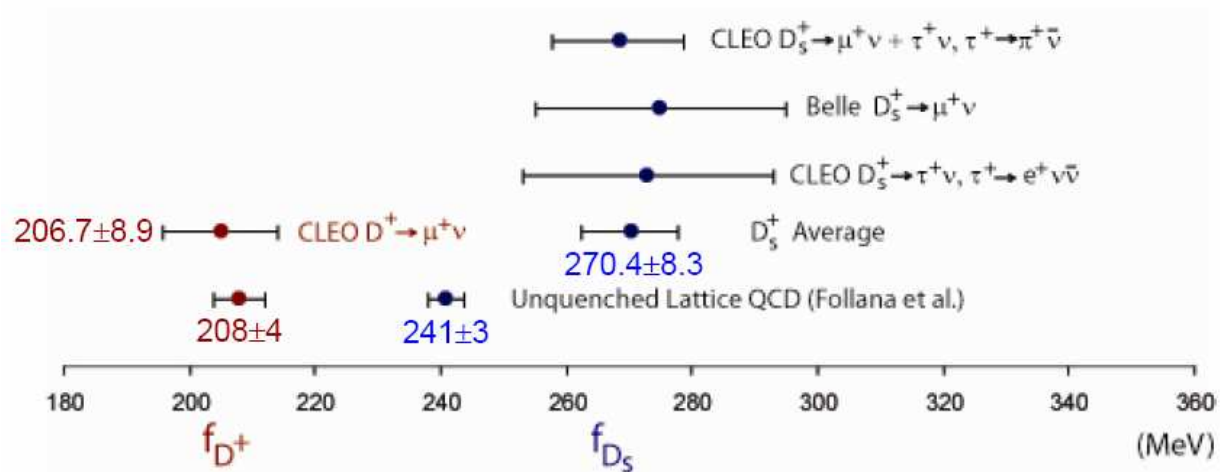
$$\text{BR}(D_s \rightarrow \mu^+ \nu) = (6.13 \pm 0.44 \pm 0.20) \times 10^{-3}$$

$$f_{D_s} = (268.2 \pm 9.6 \pm 4.4) \text{ MeV}$$

(result with τ/μ fixed
at SM expectation)

D^+ and D_s Decay Constants: the Global Picture

f_D agrees well with lattice QCD; f_{D_s} measurements internally consistent, but more than 3.5 sigma away from lattice QCD !



Is there something wrong with calculation (implications for mixing side) or is it new physics (charged Higgs, leptoquarks... arXiv:0803.1898 ; arXiv:0803.0512)?

Final D_s results from CLEO-c expected soon with full data sample

New Physics Searches in the D system

-
- D^0 - \bar{D}^0 oscillations (in brief – see next talk)
 - The search for CP violation
 - Rare decays

$D^0-\bar{D}^0$ Mixing

Short-distance

$D^0-\bar{D}^0$ transitions have two observables:

$$x = \frac{\Delta M}{\Gamma}, \quad y = \frac{\Delta\Gamma}{2\Gamma}$$

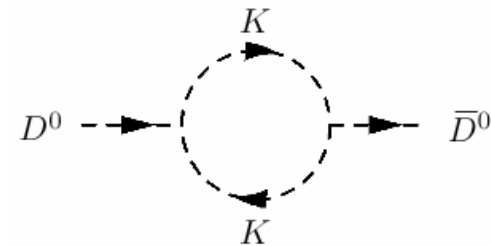


Boxes and loops in charm transitions involve down-type quarks – this gives charm system unique new physics sensitivity.

SM calculations based on box diagrams alone gives $x \sim 10^{-5}$, $y \sim 10^{-7}$ [Falk et al. PRD 65 (2002) 054034]

y should be dominated by long-distance effects, and is generally considered to be immune to NP (but not always, ie. [Petrov & Yegiyany PRD 77 (2008) 034018]). So $x \gg y$ would point to NP!

Long-distance

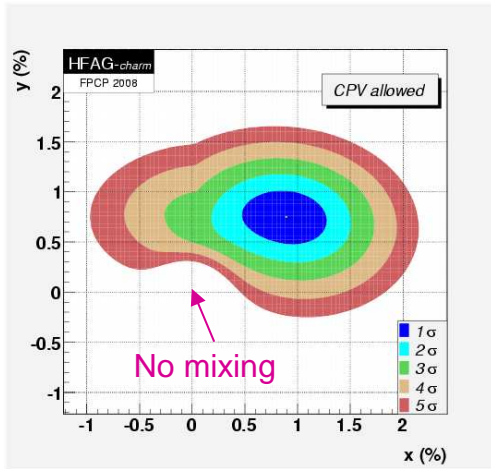
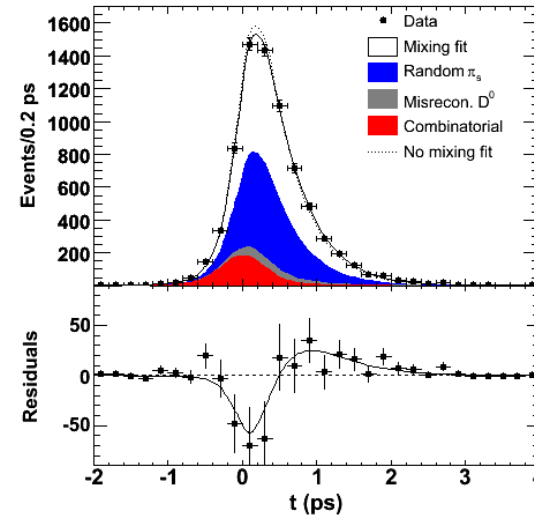


D⁰- \bar{D}^0 Mixing: Observation

Numerous recent, exciting results on charm mixing (see Joerg Marks' talk).

The most interesting....

- 'Wrong sign' $K\pi$ (x^2, y')
 - BELLE PRL 96 (2006) 151801
 - BaBar PRL 98 (2007) 211802
 - CDF PRL 100 (2008) 121802
- Eigenstate lifetime analyses: y_{CP}
 - BaBar arXiv:0712.2249
 - BELLE PRL 98 (2007) 211803
- $K_S \pi^+ \pi^-$ Dalitz analyses: x, y
 - BELLE PRL 99 (2007) 131803



No doubt now that mixing has been seen...

$$x = 0.89 \pm_{0.27}^{0.26} \%$$

$$y = 0.75 \pm_{0.18}^{0.17} \%$$

(HFAG May 08,
CPV allowed)

...but what does it mean?

D⁰- \bar{D}^0 Mixing: Interpretation

SM predictions for x & y have very large spread in value.

Observed values of parameters:

$$x = 0.89 \pm_{0.27}^{0.26} \%$$

$$y = 0.75 \pm_{0.18}^{0.17} \%$$

are on high side of what was expected, but are consistent.

For this reason, and since $x \sim y$, no immediate sign of new physics, but plenty of useful constraints can be derived

Golowich, Hewett, Pakvasa and Petrov, PRL 98 (2007) 181801

Model	Approximate Constraint
Fourth Generation	$ V_{ub'}V_{cb'} \cdot m_W < 0.5$ (GeV)
$Q = -1/3$ Singlet Quark	$s_2 \cdot m_S < 0.27$ (GeV)
$Q = +2/3$ Singlet Quark	$ \lambda_{uc} < 2.4 \cdot 10^{-4}$
Little Higgs	Tree: See entry for $Q = -1/3$ Singlet Quark Box: Parameter space can reach observed x_D
Generic Z'	$M_{Z'}/C > 2.2 \cdot 10^3$ TeV
Family Symmetries	$m_1/f > 1.2 \cdot 10^3$ TeV (with $m_1/m_2 = 0.5$)
Left-Right Symmetric	No constraint
Alternate Left-Right Symmetric	$M_R > 1.2$ TeV ($m_{D_1} = 0.5$ TeV) $(\Delta m/m_{D_1})/M_R > 0.4$ TeV ⁻¹ $M_{VLQ} > 55(\lambda_{PP}/0.1)$ TeV
Vector Leptoquark Bosons	No constraint
Flavor Conserving Two-Higgs-Doublet	$m_H/C > 2.4 \cdot 10^3$ TeV
Flavor Changing Neutral Higgs	$m_H/ \Delta_{uc} > 600$ GeV
FC Neutral Higgs (Cheng-Sher)	See entry for RPV SUSY
Scalar Leptoquark Bosons	$M > 100$ TeV
Higgsless	No constraint
Universal Extra Dimensions	$M/ \Delta y > (6 \cdot 10^2)$ GeV
Split Fermion	$M_1 > 3.5$ TeV
Warped Geometries	$ (\delta_{12}^u)_{LR,RL} < 3.5 \cdot 10^{-2}$ for $\tilde{m} \sim 1$ TeV
MSSM	$ (\delta_{12}^u)_{LL,RR} < .25$ for $\tilde{m} \sim 1$ TeV $\tilde{m} > 2$ TeV
SUSY Alignment	$\chi'_{12k}\chi'_{11k}/m_{\tilde{d}_{R,k}} < 1.8 \cdot 10^{-3}/100$ GeV
Supersymmetry with RPV	No constraint
Split Supersymmetry	No constraint

CP Violation in the Charm System

If D mixing discovery has not (immediately) revealed New Physics, where to look?
Answer: CP violation ! Extremely small in SM (~ only 2 generations participate).

Two possible sources of CPV (there is another – see later):

ϕ - phase between mixing and decay
 $(|q/p| - 1)$ - where $D_{+,-} = p|D^0\rangle \pm q|\bar{D}^0\rangle$ both negligibly small in SM

So CP asymmetries possible in mixing (A^m) or between mixing and decay (A^i):

$$A^m \propto -y/2 (|q/p| - |p/q|) \cos\phi$$

$$A^i \propto x/2 (|q/p| + |p/q|) \sin\phi$$

New Physics observable giving non-0 ϕ or $(|q/p|-1)$ suppressed by $x, y \sim 10^{-2}$.
So if something is seen it will be very small – but not as small as once feared !

Indirect CP Violation in Charm: Status & Prospects

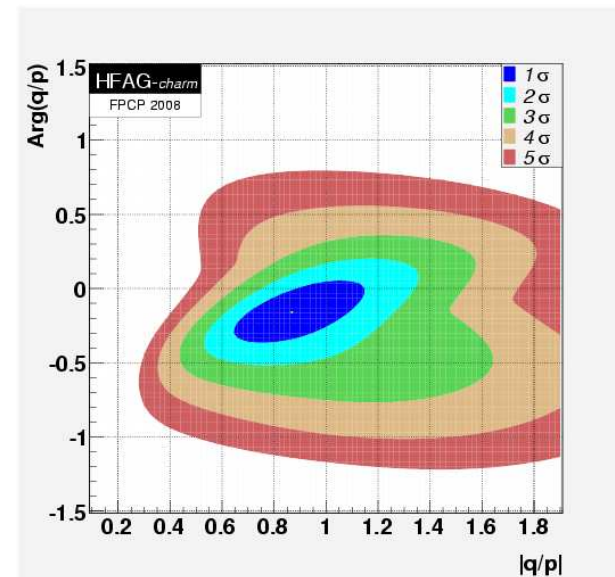
Generalising the mixing analyses to allow for CPV violation gives sensitivity to the two parameters governing CPV in mixing and mixing-decay interference.

So far no evidence of CPV, but existing limits are already quite impressive:

$$|q/p| = 0.87 \pm \begin{matrix} 0.18 \\ 0.15 \end{matrix}$$

$$\phi = -9.1 \pm \begin{matrix} 8.1 \\ 7.8 \end{matrix} \text{ degrees}$$

Higher sensitivity will come as mixing analyses improve precision – wait for LHCb (and beyond)



Many NP models expect effects here (eg. flavour alignment in SUSY)

CPV in Singly Cabibbo Suppressed D decays

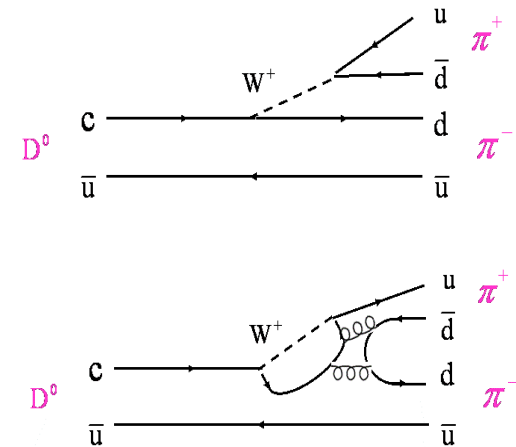
Singly Cabibbo Suppressed (SCS) decays - a win-win scenario for CPV searches

Interference between tree and Penguin can generate both direct CP asymmetries which:

- Could reach $\sim 10^{-3}$ in SM - may be observable!
- In many NP models effects of $\sim 10^{-2}$ possible (see eg. Grossman, Kagan, Nir, PRD 75 (2007) 036008)

Task therefore is to look for CPV in SCS:

- Time integrated asymmetries in CP eigenstates, eg. KK , $\pi\pi$ (involves other types of CPV, not just direct)
- Asymmetries in charged D decays, eg. $KK\pi$
- Asymmetries in final state distributions, eg. Dalitz plots and moments (may be most sensitive)



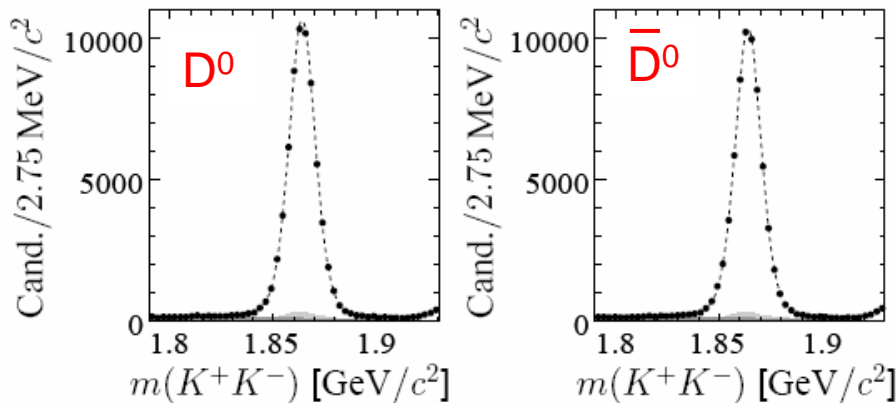
CPV searches in $D^0 \rightarrow KK, \pi\pi$

Measure asymmetry in time integrated rates:

$$A_{CP} = \frac{\Gamma(D^0 \rightarrow KK(\pi\pi)) - \Gamma(\bar{D}^0 \rightarrow KK(\pi\pi))}{\Gamma(D^0 \rightarrow KK(\pi\pi)) + \Gamma(\bar{D}^0 \rightarrow KK(\pi\pi))}$$

Distinguish D flavour from 'slow pion' charge in $D^* \rightarrow D^0\pi$

BaBar, PRD 100 (2008) 061803



386 fb⁻¹ , ~130k KK events

Spot the difference....

Use $K\pi$ events to calibrate out asymmetries in slow π reconstruction. Form CP asymmetry in bins of θ to account for EW (γ -Z) FB asymmetries.

$$A_{CP} = 0.00 \pm 0.34 \text{ (stat)} \pm 0.13 \text{ (syst)}\%$$

(dominates HFAG world average)

← Entering interesting territory !

LHCb will accumulate > 50x statistics in this mode – big improvements possible...

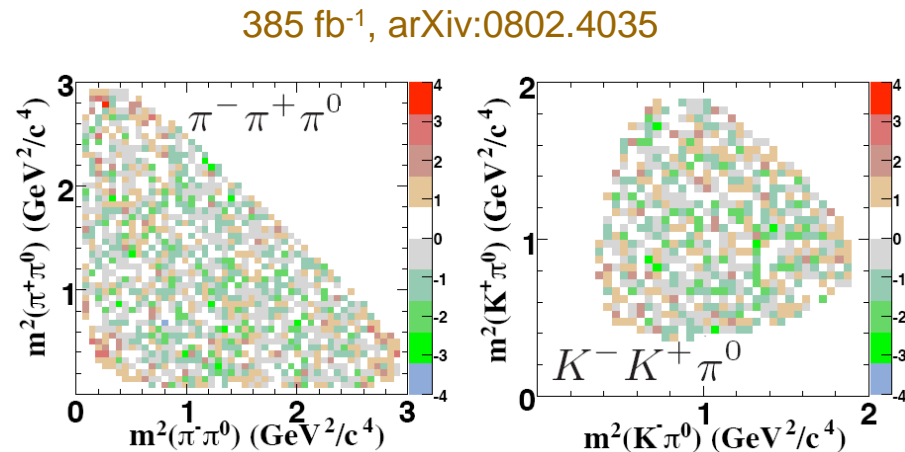
CPV Searches in Multibody ($n \geq 3$) Decays

Final state distributions in 3 and 4-body decays allow for other strategies, some with higher sensitivity than simple comparisons of integrated rates.

eg. BaBar study of $K^+K^-\pi^0, \pi^+\pi^-\pi^0$

Several complementary analyses:

- Form residuals of D^0, \bar{D}^0 w.r.t. mean in Dalitz space
- Look for difference in angular moments of D^0 & \bar{D}^0 distributions
- Compare amplitude fits of D^0 & \bar{D}^0 Dalitz spaces (model dependent)
- Look for phase space integrated asymmetry.



Other example: FOCUS T-odd moments study in $D \rightarrow KK\pi\pi$ (PLB 622 (2005) 622)

New Physics Searches with Rare D Decays

Example: $c \rightarrow u l^+ l^-$. Initially, exclusive modes look unpromising for NP searches, in contrast to B decays.

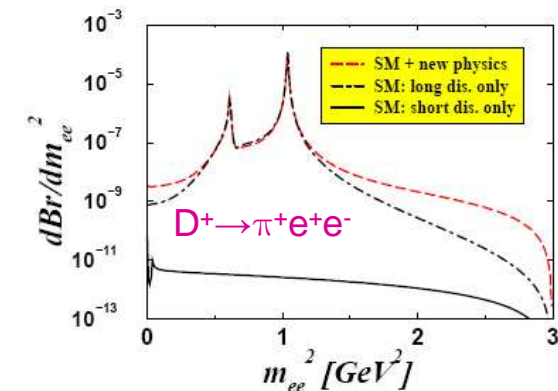
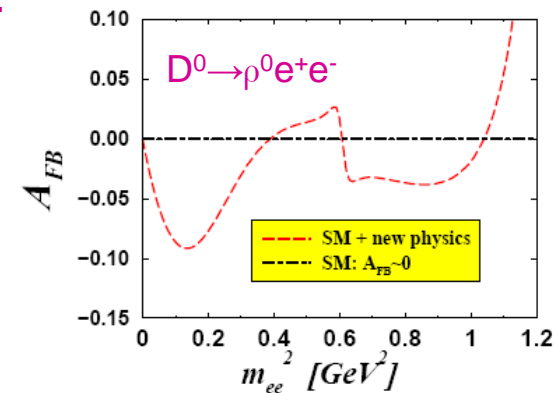
Why? Because short distance effects are swamped by long distance contributions.

Br	short distance contribution only		total rate \simeq long distance contr.	experiment
	SM	SM + NP		
$D^+ \rightarrow \pi^+ e^+ e^-$	6×10^{-12}	8×10^{-9}	1.9×10^{-6}	$< 7.4 \times 10^{-6}$
$D^+ \rightarrow \pi^+ \mu^+ \mu^-$	6×10^{-12}	8×10^{-9}	1.9×10^{-6}	$< 8.8 \times 10^{-6}$
$D^0 \rightarrow \rho^0 e^+ e^-$	negligible	5×10^{-10}	1.6×10^{-7}	$< 1.0 \times 10^{-4}$
$D^0 \rightarrow \rho^0 \mu^+ \mu^-$	negligible	5×10^{-10}	1.5×10^{-7}	$< 2.2 \times 10^{-5}$

However, differential distributions, and FB asymmetries, still have discriminating power.

And total rate can still be sizably enhanced in some cases: $D^0 \rightarrow \mu \mu \sim 10^{-13}$ in SM, can go up to 10^{-7} in R-parity violating SUSY

NP models with extra up-type quark. From arXiv:0801.1833



Conclusions

Charm has stepped back into the front-line of flavour physics

- SM description of CP violation has withstood its first attack (from the B-factories and Tevatron). Next phase of measurements, at LHCb, will have heavy reliance on what we understand about D decays.
- The D meson system is an excellent place to look for new physics (and one which is complementary to the B sector). In particular the \sim zero expected CP violation in the SM gives a near background-free environment in which to search. The observation of D-mixing gives us heart.

So, as well as being instrumental in the establishment of the SM, charm may yet play a role in its dethronement!