DISCUSSION OF POSSIBLE EVIDENCE FOR NON-LINEAR BCS RESISTANCE IN SRF CAVITY DATA TO MODEL COMPARISON



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Powerful RF cavities are now being developed for future large-scale particle accelerators from 1.E+12 1.E+1 high purity sheet niobium (Nb) superconductor, reaching peak RF surface magnetic fields of up 1.E+1 lin. BCS only, k=7.6, h=1780 lin. BCS only k=6.9/9.25, (II) w non-lin BCS-corr lin BCS -cell) -cell) w. non-lin BCS corr (C~2) Bc=180 mT k=11.2, h=3956 to 180mT. The basic model for Q-slope in SRF cavities is the thermal feedback model (TFBM). (C~3.9/2.9) Bc=180 mT h=1445/2699 the result of the exponential dependence of the Nb BCS surface resistance on temperature and 1.E+11 · · · · · · · (TESL the dependence of the RF power dissipation on the surface resistance. Most important for the Ë 5 1.E+10 1.E+1 validity of the TFBM is what surface resistance contributions (beyond BCS) are included. Here factor CEA C115 1.6K we discuss if the non-linear correction to the BCS resistance as recently proposed by Gurevich 1.E+10 before/after bake Cornell EI1-30. DESY - AC70.~1.9 H 1.53K /1.75K lity ality could be one of them, comparing TFBM calculations with experimental bulk Nb cavity data from 10 25 non-lin BCS co before/after baking H_=180mT) k=7.6, h=1780 DESY, CEA, J-Lab, Cornell and FNAL. ş ā 1.E+09 1.E+09 1 E+09 0 40 80 120 160 200 40 80 120 160 200 40 120 160 0 80 200 Peak RF field (mT) Peak RF field (mT) Peak RF field (mT The Thermal Feedback Model (TFBM) 1.E+11 Comparison of TFBM and Cavity Data The steady state heat balance equation (Eq.1) contains conduction and generation terms. The deltaw non-lin BCS-corr lin BCS only (C~2) Bc=180 m k=12.7. h=502 function in the generation term reflects the fact that the RF heating is concentrated in a very thin surface TFBM parameters for cavity Q calculation. Linear (Δ/k_BT_c, layer. The RF power dissipated per unit area in the cavity depends on the RF magnetic field amplitude Hpe DESY 2 (1.9) FNAI 1.8 291 2.6 9.9 JLAB 2.0 282 2.6 12.7 CORNELL 1.6 283 2.6 7.6 1780 1 (2) and the (temperature dependent) RF surface resistance R_s(T) as given in Eq. 2. The equation assumes T₂ (K) $\begin{array}{l} T_{0}\left(K\right)\\ G\left(\Omega\right)\\ d\left(mm\right)\\ \kappa(T_{0})\mathcal{W}/K/m\right)\\ h_{top}(T_{0})\mathcal{W}/K/m^{2})\\ R_{m}\left(n\Omega\right)\\ R_{m}\left(n\Omega\right)\\ R_{bcs,h}(T_{0})\left(n\Omega\right)\\ \Delta/k_{B}T_{c}\\ \end{array}$ A(w)) and non-linear BCS 283 2.6 6.1 1090 270 2.6 11.22 273 2.6 5.8 255 2.75 that the loss is due to the RF shielding currents only and neglects the contribution by electric surface fields. ö resistance (C(T₀,ω)), thermal 0000 1 E + 106.9 (9.3) The solution of Eq. 1 depends on the surface temperatures on both sides of conductivity $(\kappa(T_0))$ and 3956 3080 5021 956 1445 (269 3.2 (4.2) -10 (5.2) 10 40 17 (9.4) 3.6 (5) 11 (11) 5.6 (1) the niobium sheet. The temperature on the RF exposed side, T_m, drives the Kapitza conductance (h_{Kap}). 05(03) 17(105) 24 (4.3) 1.53 (1.94) 31 (20) 2.1 (1.94) 39(51 surface resistance, while the temperature on the helium side, T, drives the JLAB - 11 SC. 2.0 K 1.92 14.8 9.2 3.9 2.9 $\frac{\partial}{\partial x} \kappa(T) \frac{\partial T}{\partial x} + P_{diss}(T_m, H_c, ...) \delta(x) = 0 \qquad \left(\frac{W}{m^2}\right)^2$ 2 (2.05) 1.97 (1.93 2.09 (2.15) 1.99 (1.99 3.7 (2.5) Data out(inside) parentheses ₹ A(ω) (10⁻⁵ Ω) T_c (K) * 2 76(2 13) 25(12) 0 597(1 058) 44(17) 4 46 (2 38) Kapitza interface conductance. They can be derived exactly from the boundary are for before(after) the low 9.2 1.3 9.2 1.3 9.2 1.5 9.2 1.5 9.2 1.5 9.2 ω/2π (GHz) C (T₀,ω) conditions (Eqs. 3 & 4) for a given HRF, To and Rs (Tm,HRF,..). We used a temperature bake. * assumed 4.5 (4.5) 15(25) 2.6 (2.2) 5.2 (5.5) $\boldsymbol{P}_{diss} = \frac{1}{2} \boldsymbol{R}_{s} (\boldsymbol{T}_{m}) \boldsymbol{H}_{RF}^{2} \qquad \left(\frac{\boldsymbol{W}}{m^{2}} \right)$ 1 E+09 36(34) 39(29) computer program to calculate the exact, numerical solutions of Eqs. (3)&(4). values 80 120 160 0 40 200 The strong temperature dependence of the BCS resistance is at the core of Peak RF field (mT) thermal feedback. The increase of the surface resistance with field is the result $h_{Kap}(T_s, T_0)d(T_s - T_0) = \int_{-\infty}^{T_s} \mathbf{k}(T')dT' \qquad \left(\frac{W}{W}\right)$ The model consists of the exact solution of the TFBM equations, using the linear BCS and residual resistances measured in the cavities at low field and calculated of a feedback process during which the surface temperature increases due to material properties. Calculations were performed with and without the non-linear correction to the BCS resistance. Note that the model implementation here assumes RF heating while the RF heating increases with surface temperature. The uniform surface properties. The most important criterion the experimental data needed to satisfy for this comparison is that they needed to have as little Q slope as $\frac{1}{2}\boldsymbol{R}_{s}(\boldsymbol{T}_{m},\boldsymbol{H}_{c},...)\boldsymbol{H}_{RF}^{2} = \boldsymbol{h}_{Kap}(\boldsymbol{T}_{s},\boldsymbol{T}_{0})(\boldsymbol{T}_{s}-\boldsymbol{T}_{0}) \qquad \left(\frac{\boldsymbol{W}}{\boldsymbol{m}^{2}}\right)$ TFBM is only as good as the surface resistance and thermal parameter models possible, such as to limit as much as possible the surface resistance to the basic residual and BCS components. Most cavities were reduced size prototypes, with the only that are put in exception being the DESY AC70 (9-cell TESLA cavity). The Saclay and DESY cavities were electro-polished, the J-Lab, Cornell and FNAL cavities were BCP etched. The J-Lab cavities and the Saclay cavity C115 were also post-purified. The data obtained before and after the low temperature (-120°C, 50 hrs) bake are presented. Material Parameters Essentially all Q measurements were performed in the CW mode. Table 1 summarizes the experimental and model parameters used in the comparison. Surface Resistance lin, bCS only, k=6.1, h=1090 $R_{s,BCS}(T) = A(\boldsymbol{\omega}) \frac{T_c}{T} e^{-\boldsymbol{\omega} \frac{T_c}{T}}$ (Ω) The RF surface resistance of bulk, high purity Nb is usually defined as a sum of 1 0E+1 lin. BCS only 1 E+1 1 0E+11 -cell) k=9.9, BCS resistance (R_{s BCS}), and residual resistance (R_{res}). A phenomenological fit of the linear R_{s BCS} is given on the left ($\alpha = \Delta/k_BT_c \sim 2.0$). $\mathbf{R}_{s,BCS}(\mathbf{T},\mathbf{H}_{RF}) = \mathbf{R}_{s,BCS}(\mathbf{T},\mathbf{H}_{RF}=0) \times ...$ As recently discussed by A. Gurevich the BCS contribution increases at fields (CEBAF 1.E+10 lin. BCS only k=5.8, h=956 approaching the critical field as a result of distortions of the electronic band w. non-lin. BCS-corr $... \times 1 + C(T, \omega) \left(\frac{H_{RF}}{H}\right)$ FNAL - 3harm-3-cell (C=4.5, Bc=180mT (TESLA-: structure in the superconductor. The first critical field correction term to BCS Ē 1.0E+09 w. non-lin BCS-cor 1.0E+10 1.8 K. 3.9 GHz goes with C(T, w), a constant of order unity in Nb at ~ 2 K, that can be calculated actor (C~5.3) Bc=180 mT from material parameters w. non-lin. BCS-corr CEA C103 1.44K Quality JLAB - OCSC, 1.4 K The figure shows different model implementations of the thermal conductivity Quality (C=2.9, B_=180mT before/after bake before/after baking of polycrystalline, high purity Nb, ø consistent with experimental 1.E+09 1.0E+08 1.0E+09 60 80 100 120 data. Instead of using a full-40 80 120 160 200 0 20 40 0 40 80 120 160 200 0 Peak RF field (mT) blown model (such as Koechlin Peak RF field (mT) Peak RF field (mT) Thermal Properties and Bonin for instance) we The non linear BCS resistance contribution decreases with temperature and increases with frequency. We observed that the CEA C-103 and J-Lab-OCSC cavities, which $\kappa(T) = 0.7 e^{(1.65T - 0.1T^2)}$ used a simple fit. Note that this Koechlin-Bonin (RRR 300), n were tested at lower temperatures than the others (discussed here), show a medium-Q slope that is more pronounced than that which can be predicted even with the fit assumes a "mild" phonon phonon peak addition of the non-linear BCS resistance in the TFBM. Reduced thermal parameters cannot explain this discrepancy. Could that indicate that the increase of non-linear peak. Similarly we used a Koechlin-Bonin (RRR 300) BCS resistance is underestimated? Is it that at very low temperature the resistance is not BCS dominated anymore and some "other" process takes place? with phonon peak phenomenological fit for the $h_{Kap}(T, T_0) = 200 \cdot (T_0^{4.65})$ Solyak (300) based o Reschke/DESY data At frequencies beyond 2 GHz, the experimental data are well-described using only linear BCS resistance and residual resistance in the TFBM - see the case of the Kapitza interface conductance. $+\left(\frac{T-T_{0}}{T}\right)^{2}+0.25\left(\frac{T-T_{0}}{T}\right)^{2}$ $\overline{Km^2}$ Fermilab 3rd harmonic cavity, which operates at 3.9 GHz (In this case the TFBM with linear BCS only predicts even the exact quench field!). This discrepancy is surprising such as proposed by Mittag 10 since the high frequency regime is BCS resistance dominated! Does this indicate that the non-linear BCS resistance model as formulated here overestimates the non-0 6 for T-T₀<1.4 K.

linear BCS resistance increase at higher frequency? Or is this an indication of a more fundamental flaw of the non-linear BCS model?

Temp (K)