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# FURTHER SEARCH FOR THE TWO-PHOTON PRODUCTION OF THE GLUEBALL CANDIDATE $f_J(2220)$

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## Abstract

The CLEOII detector at the Cornell $e^+e^-$ storage ring CESR has been used to search for the two-photon production of the $f_J(2220)$ decaying into $\pi^+\pi^-$. No evidence for a signal is found in data corresponding to an integrated luminosity of $4.77\,\mathrm{fb}^{-1}$ and a 95% CL upper limit on $[\Gamma_{\gamma\gamma}\,B_{\pi^+\pi^-}]_{f_J(2220)}$ of $2.5\,\mathrm{eV}$ is set. If this result is combined with the BES Collaboration's measurement of $f_J(2220) \to \pi^+\pi^-$ in radiative $J/\psi$ decay, a 95% CL lower limit on the stickiness of the $f_J(2220)$ of 73 is obtained. If the recent CLEO result for $[\Gamma_{\gamma\gamma}\,B_{K^0_S K^0_S}]_{f_J(2220)}$ is combined with the present result, the stickiness of the $f_J(2220)$ is found to be larger than 102 at the 95% CL. These results for the stickiness (the ratio of the probabilities for two-gluon coupling and two-photon coupling) provide further support for a substantial neutral parton content in the $f_J(2220)$.

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In lowest order, the two-photon width of a resonance is proportional to the fourth power of the constituent parton charges, so a very small two-photon width is an indication of substantial neutral parton content. Within the framework of QCD, a small two-photon width implies that the resonance has substantial glueball content. A quantitative measure of the glueball content of a resonance is the ratio of the probabilities for two-gluon coupling and two-photon coupling for which the resonance's two-gluon coupling is deduced from its production rate in radiative $J/\psi$ decay.

The $f_J(2220)$ is a glueball candidate owing to its observation in radiative $J/\psi$ decay (a glue-rich environment) [1,2], its small two-photon width relative to its two-gluon width [3,4], its small total width [1,2], its similar branching fraction for non-strange and strange final states [2], and its proximity to the mass obtained in lattice calculations [5,6] for a tensor glueball. CLEO has recently [4] obtained a 95% CL upper limit on the product of the two-photon width and the $K^\circ_S K^\circ_S$ branching fraction $[\Gamma_{\gamma\gamma} B_{K^\circ_S K^\circ_S}]_{f_J(2220)}$ of 1.3 eV using the reaction $e^+e^- \to e^+e^- f_J(2220) \to e^+e^- K^\circ_S K^\circ_S$. Earlier, the ARGUS Collaboration [3] obtained a less restrictive limit based upon the $K^+K^-$ decay mode. In the present paper we report on a search for the two-photon production of the $f_J(2220)$ in the reaction $e^+e^- \to e^+e^- f_J(2220) \to e^+e^-\pi^+\pi^-$.

The CLEOII detector [7] is a general purpose detector operating at the Cornell Electron Storage Ring CESR [8]. It provides charged particle tracking, precision electromagnetic calorimetry, charged particle identification, and muon detection. Charged particle detection over 95% of the solid angle is provided by three concentric drift chambers in a magnetic field of 1.5 T giving a momentum resolution $\sigma_p/p = 0.5\%$ at $p = 1$ GeV. The drift chambers are surrounded by a time of flight system and a CsI electromagnetic calorimeter. A superconducting coil and muon detectors surround the calorimeter. Two-prong events are recorded with three redundant triggers. The results in this paper are based upon an integrated luminosity of 4.77 fb$^{-1}$ with CESR operating at a center-of-mass energy of approximately 10.6 GeV.

The $f_J(2220)$ is searched for in the two-photon reaction $e^+e^- \to e^+e^- f_J(2220) \to e^+e^-\pi^+\pi^-$ in the untagged mode in which the outgoing $e^+$ and $e^-$ are undetected. Events are selected that have exactly two tracks of opposite charge whose vector sum of momenta transverse to the beam has a magnitude less than 0.5 GeV. The total energy of the event is required to be less than 6.0 GeV and the energy in the calorimeter not associated with either track must be less than 0.5 GeV.

Two-photon produced final states of charged particle pairs are selected (and backgrounds from Bhabha scattering, muon pair production, and cosmic rays are suppressed) by requiring that the acolinearity of the two tracks is greater than 0.1. In addition, the acoplanarity is required to be less than 0.05. Here acolinearity is the deviation from colinearity in three dimensions while acoplanarity is the deviation from colinearity in the plane transverse to the beams. These last two requirements are effective because the two-photon center-of-mass generally moves rapidly and at a small angle with respect to the beams.

Events are vetoed if either track is identified as an electron or muon. If $E/p$, the ratio of a track's energy deposition in the calorimeter and its momentum measured in the drift chambers, is in the range $0.85 - 1.10$, the track is identified as an electron. Muons are identified by the muon detectors. Events must have satisfied at least one of the two-prong triggers.

The event simulation uses the BGMS [9] formalism with the transverse-transverse term (appropriate for untagged two-photon reactions) for the event generation and GEANT [10] for the detector simulation. Photon form-factors based upon vector-meson dominance with a mass $m_V = 768.5\,\mathrm{MeV}$ are used. We take the spin of the $f_J(2220)$ to be 2. The detection efficiencies for helicity 0 and 2 are found to be 13.1% and 26.9% respectively. We use a ratio [11] of helicity 0 and helicity 2 of 1:6, giving an efficiency of 24.9%. When the mass $m_V$ in the photon form-factors is varied from 768.5 MeV to $\infty$ (corresponding to a form-factor equal to 1) the cross-section increased by 29.8% while the efficiency dropped by 18.9% and their product increased by 5.5%. A 2.8% systematic uncertainty is assigned to the product of the cross-section and efficiency.

The dominant source of systematic uncertainty is the trigger efficiency. It is estimated to be 13% from the observed variation of the event yield as a function of the azimuthal angle of each of the two tracks. Data and simulation are compared to determine smaller systematic uncertainties of 2% per track from track reconstruction efficiency, 3% from the requirement on the energy deposition in the calorimeter, 3% from the transverse momentum requirement, 2% each from the acolinearity and acoplanarity requirements, 5% from the $E/p$ requirement, and 4% from the muon veto. The total systematic uncertainty is the sum in quadrature of the above sources and is 15%.

A pion-pair invariant mass distribution is constructed using all events that pass the selection criteria and assuming that both particles are pions. A plot of $m_{\pi^+\pi^-}$ in the mass region relevant for the $f_J(2220)$ is shown as the data points with statistical error bars in Fig.1. There is no evidence of an enhancement near the mass of the $f_J(2220)$. The mass distribution is fit with the sum of a signal and a background assuming that there is no interference between the two. The signal shape is represented by a Breit-Wigner with a mean of 2231 MeV [12] and a width of 23 MeV [12] convolved with the detector resolution of 12 MeV and is shown as the hatched histogram in Fig.1. The background is represented by a third order polynomial that is fit to the mass region $2000-2500\,\mathrm{MeV}$ excluding the region $2200-2268\,\mathrm{MeV}$. The fit gives a signal of $-103 \pm 77$ events with a $\chi^2 = 35.6$ for 36 degrees of freedom.

An upper limit is obtained by only allowing for a positive number of signal events, $N$. Given that $m_{f_J(2220)} = 2231.1 \pm 2.5\,\mathrm{MeV}$ [12] and $\Gamma_{f_J(2220)} = 23^{+8}_{-7}\,\mathrm{MeV}$ [12], likelihood functions for $N$ are obtained for a range of the resonance mass and width, spanning $\pm 2.5\sigma$ in each. These functions are then weighted with Gaussian probabilities for the mass and width to obtain a final likelihood function $L_N$. The product of the two photon partial width and charged di-pion branching fraction, $\Gamma_{\gamma\gamma}\,B_{\pi^+\pi^-}$, is given by the product of $N$ and $P$. Here $P$ is the partial width used in the simulation divided by the product of luminosity, cross-section and efficiency; $P$ is assumed to be Gaussian distributed. The likelihood function, $L_{\Gamma B}$ is then obtained by numerical integration in the two-dimensional space of $N$ and $P$. From $L_{\Gamma B}$ a 95% CL upper limit of 2.5 eV for $\Gamma_{\gamma\gamma}\,B_{\pi^+\pi^-}$ is obtained. The solid line in the main portion of Fig.1 is the sum of the fit to the background and a signal that corresponds to this upper limit. The mass region $2150-2310\,\mathrm{MeV}$ is shown enlarged in the inset in Fig.1 with the two curves representing the background fit with and without this level of signal added.

The upper limit can be specified without the assumption of a 1:6 ratio for helicity 0 and 2 as $(0.53\Gamma^{2,0}_{\gamma\gamma} + 1.08\Gamma^{2,2}_{\gamma\gamma})B_{\pi^+\pi^-} < 2.5\,\mathrm{eV}$ at 95% CL. The superscripts indicate spin and

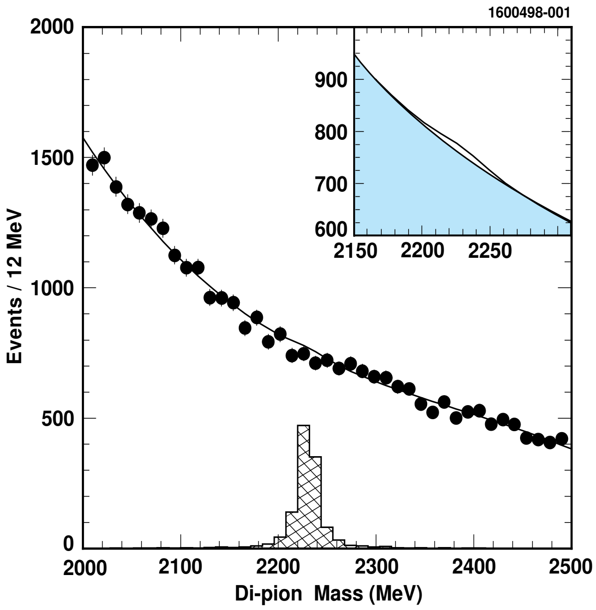


FIG. 1. The $\pi^+\pi^-$ invariant mass distribution for the data in the region of the $f_J(2220)$. The hatched histogram is the expected signal shape with arbitrary normalization. The solid curve is the sum of a fit to the background and a signal corresponding to the 95% CL upper limit on $\Gamma_{\gamma\gamma}B_{\pi^+\pi^-}$ of 2.5 eV. In the insert the two curves are the background fit with and without this level of signal added.

helicity. The ratio of the coefficients is equal to the ratio of the efficiencies for helicity 0 and $\pm 2$ while the overall normalization is determined by the result given above.

The upper limit on $\Gamma_{\gamma\gamma}\, B_{\pi^+\pi^-}$ can be interpreted in terms of the stickiness $S$ [13]. Stickiness is the ratio of the probabilities for two-gluon and two-photon coupling of a resonance, which in the present case can be written as ($f_J$ denotes $f_J(2220)$):

$$S_{f_J} = \frac{|\langle f_J|gg\rangle|^2}{|\langle f_J|\gamma\gamma\rangle|^2} = C_\ell \left(\frac{m_{f_J}}{k_\gamma}\right)^{2\ell+1} \frac{\Gamma_{J/\psi} B(J/\psi \to \gamma f_J) B(f_J \to \pi^+\pi^-)}{\Gamma(f_J \to \gamma\gamma) B(f_J \to \pi^+\pi^-)} \tag{1}$$

The parameter $k_\gamma$ is the energy of the photon produced in the radiative $J/\psi$ decay as calculated in the $J/\psi$ rest frame, and $\Gamma_{J/\psi}$ is the total width of the $J/\psi$. The factor with $2\ell+1$ in the exponent removes the trivial phase space dependence of the stickiness upon the $f_J$ mass. The quantum number $\ell$ is the relative angular momentum between the two gluons or photons, with $\ell = 0$ for $J = 2$. $C_0 = 20.5$ is a normalization factor chosen such that the stickiness is normalized to unity for the $f_2(1270)$. The BES result [2] and $J/\psi$ properties from the Particle Data Group [12] are combined with our result to obtain a likelihood distribution for the stickiness of the $f_J$ via a Monte Carlo technique. In this procedure the $L_{\Gamma B}$ obtained previously was used and all other uncertainties were taken to be Gaussian distributed. A lower limit of $S_{f_J} > 73$ is found at 95% CL.

This lower limit and the one obtained in the $K^\circ_S K^\circ_S$ channel [4] can be merged, again using a Monte Carlo procedure, to obtain a combined lower limit [14] on the stickiness of $S_{f_J} > 102$, also at 95% CL. This result can be compared with the stickiness of the $f_2'(1525)$, a resonance thought to be predominantly an $s\bar{s}$ bound state. Using the properties of the $f_2'(1525)$ from the Particle Data Group [12] a stickiness $S_{f_2'} = 14.7 \pm 3.9$ is found, considerably smaller than the lower limit $S_{f_J} > 102$. A linear superposition of $|q\bar{q}>$ states can be constructed such that the two-photon width is negligible; the coefficients would have to take on very specific values so this possibility is considered unlikely. The large lower limits on the stickiness of the $f_J(2220)$ are therefore an indication of substantial neutral parton or glueball content.

In this Letter a restrictive 95% CL upper limit $[\Gamma_{\gamma\gamma} B_{\pi^+\pi^-}]_{f_J(2220)} < 2.5\,\mathrm{eV}$ is presented. Using the BES Collaboration's result for $f_J(2220) \to \pi^+\pi^-$ in radiative $J/\psi$ decay, this upper limit leads to a lower limit on its stickiness $S_{f_J(2220)} > 73$ at 95% CL. When these results are combined with an earlier CLEO result [4], a lower limit on the stickiness of 102 at 95% CL is obtained. This large value is difficult to understand if the valence partons of the $f_J(2220)$ are quarks and antiquarks only; therefore, the $f_J(2220)$ is likely to have a substantial neutral parton or glueball content.

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[14] In this report and in Ref. [4] the uncertainty on $C_0$ is not included because the stickiness could have been normalized to any $J = 2$ resonance that is thought to have light quarks as valence partons, leading to different values for the stickiness. If that uncertainty is included, the 95% CL lower limits on the stickiness of the $f_J(2220)$ are 69 and 67 for the $K_S^\circ K_S^\circ$ and $\pi^+\pi^-$ channels respectively (instead of 76 and 73) while the combined 95% CL lower limit is 94 (instead of 102).